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# Thermodynamic evaluation of the Fe-Th system and glass forming ability in the Fe-Th system

Kyung Jong Lee  
*Iowa State University*

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forming ability in the Fe-Th system**

Lee, Kyung Jong, Ph.D.

Iowa State University, 1987

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300 N. Zeeb Rd.  
Ann Arbor, MI 48106



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**Thermodynamic evaluation of the Fe-Th system and  
glass forming ability in the Fe-Th system**

**by**

**Kyung Jong Lee**

**A Dissertation Submitted to the  
Graduate Faculty in Partial Fulfillment of the  
Requirements for the Degree of  
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**Department: Materials Science and Engineering  
Major: Metallurgy**

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**~~In~~ Charge of Major Work**

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**For the Graduate College**

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**1987**



## TABLE OF CONTENTS

	Page
ABSTRACT	1
OVERALL INTRODUCTION	2
PART I. THERMODYNAMIC EVALUATION OF THE FE-TH SYSTEM	4
INTRODUCTION	5
THEORY	6
Evaluation of Fe-Th system	6
Estimation of thermodynamic properties	7
Use of invariant reactions	11
RESULTS AND DISCUSSION	14
SUMMARY AND CONCLUSIONS	19
REFERENCES	37
PART II. GLASS FORMING ABILITY IN THE FE-TH SYSTEM	38
INTRODUCTION	39
THEORY	40
RESULTS AND DISCUSSION	45
T <sub>0</sub> curves determination	45
TTT curves determination	45
CONCLUSIONS	48
REFERENCES	74
OVERALL SUMMARY	75
LITERATURE CITED	77
ACKNOWLEDGMENTS	78

## ABSTRACT

A calculation of the Fe-Th phase relationships was undertaken in order to develop a self-consistent description of the system. A subregular solution model was used to describe the solution behavior. Intermediate phases were treated as composition invariant. Due to the limited availability of experimental thermodynamic data, a multi-step analysis was employed to develop a better description of the thermodynamic properties in the system.

As a quantitative guide to the glass forming ability (GFA) in the Fe-Th system,  $T_0$  and TTT curves were evaluated from the previously derived thermodynamic description of the system. From the TTT curves, critical cooling rates ( $R_c$ ) to avoid crystallization were determined. Since the Fe-Th system has several solid phases, the relative stability of each phase in the metastable region is discussed. Kinetic parameters such as viscosity ( $\eta$ ), glass transition temperature ( $T_g$ ) and interfacial energy ( $\sigma$ ) were estimated since experimental values were not available.

This analysis indicates that thermodynamic data alone can not be relied upon to predict the phase boundary behavior. Both thermodynamic data and phase diagram data should be used.

## OVERALL INTRODUCTION

A calculation of the Fe-Th phase relationships was undertaken in order to develop a self-consistent description of the system. A subregular solution model was used to describe the solution behavior. Intermediate phases were treated as composition invariant. Due to the limited availability of experimental thermodynamic data, a multi-step analysis was employed to develop a better description of the thermodynamic properties in the system.

First, interaction parameters for the solid and liquid solutions were evaluated by the method of Miedema et al. [1977, 1983] and Niessen et al. [1983]. Second and third, interaction parameters were calculated with invariant reactions from the phase diagram. At each step, the calculated phase diagram was compared with the evaluated one [Kubaschewski 1982].

A glass is defined as a rigid body without long range atomic order. Glass formation is controlled by a combination of thermodynamic factors, kinetic factors and processing operations [Polk and Giessen 1978]. As a quantitative guide to the glass forming ability (GFA) in the Fe-Th system,  $T_0$  and TTT curves were evaluated from the previously derived thermodynamic description of the system.  $T_0$  curves, which are the loci of points defining  $\Delta_M G$  (solid

phase) =  $\Delta_M G$  (liquid phase), are the thermodynamic limits for the partitionless crystallization of equilibrium phases. TTT curves, which result from a combined kinetic and thermodynamic approach, are time - temperature - transformation curves.

From the TTT curves, critical cooling rates ( $R_c$ ) to avoid crystallization were determined. Since the Fe-Th system has several solid phases, the relative stability of each phase in the metastable region is discussed. Kinetic parameters such as viscosity ( $\eta$ ), glass transition temperature ( $T_g$ ) and interfacial energy ( $\sigma$ ) were estimated since experimental values were not available.

PART I. THERMODYNAMIC EVALUATION OF THE FE-TH SYSTEM

## INTRODUCTION

A calculation of the Fe-Th phase relationships was undertaken in order to develop a self-consistent description of the system. A subregular solution model was used to describe the solution behavior. Intermediate phases were treated as composition invariant. Due to the limited availability of experimental thermodynamic data, a multi-step analysis was employed to develop a better description of the thermodynamic properties in the system.

First, interaction parameters for the solid and liquid solutions were evaluated by the method of Miedema et al. [1977, 1983] and Niessen et al. [1983]. Second and third, interaction parameters were calculated with invariant reactions from the phase diagram. At each step, the calculated phase diagram was compared with the evaluated one [Kubaschewski 1982].

## THEORY

Evaluation of the Fe-Th system

A phase diagram for the Fe-Th system was initially proposed by Thomson [1965, 1966] who studied the system over its entire range of composition at temperatures between 800°C and the liquidus. Later investigations of the crystal structures of intermediate phases were made for  $\text{Fe}_{17}\text{Th}_2$  by Johnson et al. [1969], for  $\text{Fe}_3\text{Th}$  by Smith and Hansen [1965], for  $\text{Fe}_3\text{Th}_7$  by Matthias et al. [1961] and for  $\text{Fe}_7\text{Th}_2$  by Buschow and van der Goot [1971].

An evaluation of the composite data was made by Kubaschewski [1982] who proposed the diagram of Figure 1. The crystallography of the Th-rich and Fe-rich terminal solid solution phases and the intermediate phases has been summarized in Table 1. Lattice stabilities which are Gibbs energies of transformation as well as melting and transformation temperatures of Th and Fe were evaluated by Chase [1983] and are given in Table 2a. Gibbs formation energies of the intermediate phases in the system, shown in Table 2b, were determined by Skelton et al. [1973] at temperatures between 655°C and 891°C by an emf method. Metallographic studies by Thomson [1966] showed the solubility of Fe in Th at 950°C to be less than 2 at %. On the other side of the system, no data for the solubility of

Th in Fe have been found nor have liquidus lines been determined. Invariant reactions are summarized in Table 3.

### Estimation of thermodynamic properties

Gibbs energies of mixing for terminal solid solutions of Fe in Th and Th in Fe can be expressed in the forms:

$$\Delta_M G = G_{Fe} (1-X_{Th}) + G_{Th} X_{Th} + \Delta_M G^{id} + \Delta_M G^{xs} \quad (1)$$

where:

$$\Delta_M G^{id} = RT ((1-X_{Th}) \ln(1-X_{Th}) + X_{Th} \ln X_{Th}) \quad (2)$$

$$\Delta_M G^{xs} = X_{Th} (1-X_{Th}) (A + B (2X_{Th} - 1)) \quad (3)$$

$X_{Th}$ : atomic fraction of Th

In equation (1), the first and second terms represent lattice stabilities for Fe and Th, the third term represents the ideal mixing for Fe-Th alloys, and the last term represents excess Gibbs energy of mixing.

The Gibbs energy of formation for an intermediate phase is of the form:

$$\Delta_f G = \Delta_f H - T \Delta_f S \quad (4)$$

where  $\Delta_f H$  is the heat of formation and  $\Delta_f S$  is the entropy of formation. To calculate the Gibbs energies which should be in agreement with proposed diagram and the available



thermodynamic measurements, a multi-step analysis has been used.

As the first step, the approach is that of Miedema [Boom et al. 1976a, 1976b] who described the heat of formation of an intermetallic compound and the heat of mixing of liquid and solid solutions with a phenomenological model. The suggested heat of mixing of a liquid solution and heat of formation of an intermetallic compound, e.g., AB alloy is

$$\Delta_M H \text{ or } \Delta_F H = 2f(C) \frac{(C_A V_A^{2/3} + C_B V_B^{2/3})}{(N_{WS}^A)^{-1/3} + (N_{WS}^B)^{-1/3}} * \frac{(-P(\Delta\phi^*)^2 + Q(\Delta N_{WS}^{1/3})^2 - R)}{(5)}$$

where:

$$f(C) = C_A^S C_B^S \quad \text{for a liquid solution}$$

$$= C_A^S C_B^S (1 + 8(C_A^S C_B^S)^2) \quad \text{for an intermetallic compound}$$

$$C_A^S = \frac{C_A V_A^{2/3}}{(C_A V_A^{2/3} + C_B V_B^{2/3})}$$

C: atomic concentration (# of atoms/ cm<sup>3</sup>)

N<sub>WS</sub>: density of electrons at the boundary  
between dissimilar atomic cells

φ\*: chemical potential for electrons

V<sub>i</sub>: molar volume of the subscribed element

P, Q and R: the coefficients dependent on types of

## binary alloy systems

The first term takes cognizance of the chemical potential difference between the two atomic species and represents an energy reduction of the heat of mixing due to electron transfer to the more electropositive species. The second term is the positive energy cost of eliminating the electronic density discontinuity at the boundaries of Wigner-Seitz cells. The last term is an energy reduction attributable to hybridization of d-p electron bands.

In alloying, there is a volume change due to the charge transfer. The molar volume change is

$$V_A^{2/3} (\text{alloy}) = V_A^{2/3} (\text{pure}) (1 + \text{acf } C_B^S (\phi_A^* - \phi_B^*)) \quad (6)$$

where acf is a constant that has been empirically determined from experimental data to be 0.14, 0.10, 0.07 and 0.04 for monovalent, divalent, trivalent and higher valent metals. To calculate the heat of formation, the volume of alloy per mole of atoms rather than that of either element was used.

For a solid solution, the heat of alloy formation is

$$\Delta_M H(\text{solid solution}) = \Delta_M H(\text{liquid solution}) + \Delta_M H(\text{elastic energy}) + \Delta_M H(\text{structure energy}) \quad (7)$$

Here the second term arising from elastic mismatch can be estimated [Eshelby 1954, Friedel 1954] from elastic continuum theory. Thus,  $\Delta_M H$  (elastic energy) in A-B alloy may be expressed as follows:

$$\Delta_M H(\text{elastic energy}) = X_A X_B (X_A \Delta_{M A \text{ in } B}^{\text{elastic}} + X_B \Delta_{M B \text{ in } A}^{\text{elastic}}) \quad (8)$$

where

$$\Delta_{M A \text{ in } B}^{\text{elastic}} = 2\mu_B \frac{3K_A}{3K_A + 4\mu_B} \frac{(\Delta V_{AB})^2}{3V_A} \quad (9)$$

where  $\mu_B$  is the shear modulus of the B element,  $K_A$  is the bulk modulus of the A element and  $\Delta V_{AB}$  is the difference in molar volume of the two metals. The third term in equation (7) accounts for the preference for a certain crystal structure in transition metal alloy systems. This term was summarized by Miedema using theoretical estimates for the difference in energy for pure transition metals with the average number of conduction electrons per atom ( $Z$ ) [Pettifor 1977] and experimental information about solid solubilities and enthalpy differences between crystal structures.  $\Delta_M H$  (structure energy) can be expressed:

$$\Delta_M H(\text{structure}) = X_A X_B (X_A \Delta_{M A \text{ in } B}^{\text{structure}} + X_B \Delta_{M B \text{ in } A}^{\text{structure}}) \quad (10)$$

where

$$\Delta_M H_A^{\text{structure in B}} = -E_A(\phi_A) + E_B(\phi_B) + (Z_A - Z_B) \left( \frac{dE_B(\phi_B)}{dZ} \right) \quad (11)$$

where A and B are elements,  $\phi_i$  is the phase and  $E_i$  is the structure dependent energy of the subscripted element.  $E$  and  $dE/dZ$ , shown in Table 4 and Table 5, were summarized by Miedema and Niessen [1983].

#### Use of invariant reactions

As the second and third step, invariant reactions were used to optimize the interaction parameters of solution phases and Gibbs energies of formation for intermediate phases.

For example, from the 875°C temperature of the eutectic reaction (Liquid (55 at.% Th)  $\rightarrow$  Fe<sub>3</sub>Th + Fe<sub>3</sub>Th<sub>7</sub>) and the equilibrium requirement of equality of the partial molar Gibbs energies, which are,

$$\Delta_M \bar{G}_i^L = \Delta_f \bar{G}_i^{\text{Fe}_3\text{Th}} = \Delta_f \bar{G}_i^{\text{Fe}_3\text{Th}_7} \quad (i: \text{ Fe or Th}) \quad (12)$$

the following two equations were derived.

$$\Delta_M \bar{G}^L (X_{\text{Th}}=0.55, 1148\text{K}) = \left( \frac{d\Delta_M \bar{G}^L}{dX_{\text{Th}}} \right) (X_{\text{Th}}=0.55, 1148\text{K}) \\ (0.55-0.25) + \Delta_f \bar{G}^{\text{Fe}_3\text{Th}_7} (1148\text{K}) \quad (13)$$

$$\frac{d\Delta_M^{G^L}}{dX_{Th}} (X_{Th}=0.55, 1148K) = \frac{\Delta_f^{G^{Fe_3Th}}(1148K) - \Delta_f^{G^{Fe_3Th_7}}(1148K)}{0.25 - 0.7} \quad (14)$$

With the following equations for  $\Delta_f^{G^{Fe_3Th_7}}(T)$ ,  $\Delta_f^{G^{Fe_3Th}}(T)$ ,  $\Delta_M^{G^L}(T, X_{Th})$  and  $d\Delta_M^{G^L}/dX_{Th}(T, X_{Th})$ , interaction parameters (A and B) for the liquid phase can be evaluated.

$$\Delta_M^{G^{Fe_3Th_7}}(T) = -5021 - 1.67T \quad (15)$$

$$\Delta_M^{G^{Fe_3Th}}(T) = -24769 + 12.55T \quad (16)$$

$$\begin{aligned} \Delta_M^{G^L}(T, X_{Th}) = & (17406 - 9.029T)X_{Th} + (15540 - 8.8976T)(1 - X_{Th}) \\ & + RT(X_{Th} \ln X_{Th} + (1 - X_{Th}) \ln(1 - X_{Th})) \\ & + X_{Th}(1 - X_{Th})(A + B(2X_{Th} - 1)) \end{aligned} \quad (17)$$

$$\begin{aligned} (d\Delta_M^{G^L}/dX_{Th})(T, X_{Th}) = & (17406 - 9.029T) - (15540 - 8.8976T) \\ & + RT(\ln X_{Th} - \ln(1 - X_{Th})) + 2B X_{Th}(1 - X_{Th}) \\ & + (1 - 2X_{Th})(A + B(2X_{Th} - 1)) \end{aligned} \quad (18)$$

This can be accomplished by substituting equations (15)-(18) into (13) and (14) with  $T=1148$  K and  $X_{Th}=0.55$  to generate  $A = -31604$  J/mol and  $B = 1684$  J/mol.

In similar manner values from invariant reactions in Table 3 were used to optimize the interaction parameters of other solution phases and unavailable high temperature

Gibbs formation energies of intermediate phases after initially using Miedema approach to develop first approximations.

## RESULTS AND DISCUSSION

As the first step  $\Delta_M G^{XS}$  values for liquid and solid solution phases were calculated by the Miedema approach with constraints from the known invariant reactions. Throughout the calculation,  $\alpha\text{Fe}(\text{bcc})$  and  $\alpha\text{Th}(\text{fcc})$  were used as standard states.  $\Delta_M G^{XS}$  was expressed with a subregular solution model, which is,

$$\Delta_M G^{XS} = X_{\text{Th}}(1-X_{\text{Th}}) (A + B (2X_{\text{Th}}-1)) \quad (19)$$

The data in Table 12 were substituted into equation (5) to calculate the heat of mixing of a liquid solution. The values of P, Q and R in Table 12 are empirical constants that were determined by Miedema for transition metal alloy systems [Boom et al. 1976a, 1976b].

The constants in equation (5) for the liquid solution were evaluated as  $A = -42739 \text{ J/mol}$  and  $B = 11756 \text{ J/mol}$ . For the solid solution phases, the structure dependent energy and the elastic strain energy terms were calculated by substituting data from Table 13 into equations (7), (9) and (11).

This produced the following values for the excess Gibbs energies of phase formation.

$$\Delta_M G_{bcc}^{xs}(\text{structure E}) = X_{Th}(1-X_{Th}) \{-36000 + 20000(2X_{Th} - 1)\} \quad (20)$$

$$\Delta_M G_{fcc}^{xs}(\text{structure E}) = X_{Th}(1-X_{Th}) \{32000 - 20000(2X_{Th} - 1)\} \quad (21)$$

$$\Delta_M G^{xs}(\text{strain E}) = X_{Th}(1-X_{Th}) \{246000 - 99000(2X_{Th} - 1)\} \quad (22)$$

Total excess Gibbs energies for fcc and bcc phases were

$$\Delta_M G_{bcc}^{xs} = X_{Th}(1-X_{Th}) \{ 167000 - 67000(2X_{Th} - 1) \} \quad (23)$$

$$\Delta_M G_{fcc}^{xs} = X_{Th}(1-X_{Th}) \{ 235000 - 107000(2X_{Th} - 1) \} \quad (24)$$

For the intermediate phases, experimentally measured heats of formation were used. Derived Gibbs energies of all phases were summarized in Table 6. A calculated phase diagram from Table 6 was shown in Figure 2.

Calculated liquidus lines at the Th-rich side were in good agreement with the evaluate ones (Figure 1). However at the Fe-rich side, experimental liquidus lines were depressed by 100 - 200°C compared to calculated ones. This resulted from the fact that the stability of  $Fe_{17}Th_2$  was too low with respect to the stability of liquid phase to show a eutectic reaction ( $L \rightarrow \delta Fe + Fe_{17}Th_2$ ).



Moreover, compound phases showed low stabilities such that several peritectic reactions occurred at temperatures lower than expected.

In a second step, recalculation of parameters was carried out to optimize the Gibbs energies using invariant reactions and equilibrium conditions. Calculated results are listed in Table 7. Equations (4) and (5) from Table 7 were used to evaluate interaction parameters (A and B) for the formation of liquid solution. This was possible because  $\Delta_f G^{\text{Fe}_3\text{Th}}$  (1148 K) and  $\Delta_f G^{\text{Fe}_3\text{Th}_7}$  (1235 K) were experimentally available. The bcc and fcc solid phases were treated as regular solutions and their interaction parameters (C and D) were generated from equations (7) and (8) of Table 7. The entropy of formation of  $\text{Fe}_{17}\text{Th}_2$  ( $\Delta_f S^{\text{Fe}_{17}\text{Th}_2}$ ) was taken from equation (9) in Table 7. The optimization was done by linear regression. Resulting values for A, B, C, D and  $\Delta_f S^{\text{Fe}_{17}\text{Th}_2}$  were -32000 J/mol, 2000J/mol, 30000J/mol, 30000 J/mol and -2.274 J/mol deg. Calculated parameters in the second step are given in Table 8. The phase diagram is shown in Figure 3. The diagram is in good agreement with the evaluated one except that the peritectic reactions involving  $\text{Fe}_5\text{Th}$  and  $\text{Fe}_3\text{Th}$  occurred at lower temperature than expected and that  $\text{Fe}_7\text{Th}_2$  phase was disappeared in the calculated diagram.

The last step was the evaluation of the Gibbs energies

of formation of the intermediate phases. Since the experimental measurements of Gibbs energies of formation for intermediate phases were done at temperatures between 655 and 891°C which is well below the temperatures of most of the invariant reactions in the system, and because the temperature dependence of those Gibbs energies are more uncertain than the magnitude, extrapolation upward into the temperature regimes of the invariant reactions has a large inherent uncertainty. Values for  $\Delta_f H$  and  $\Delta_f S$  of the intermediate phases were therefore modified to retain consistency with the evaluated interaction parameters of the liquid and solid solution phases while remaining consistent with  $\Delta_f G$ 's in the temperature range of measurement. Equations (1) to (6) and (9) to (11) in Table 7 were used to calculate the Gibbs energies of formation for the intermediate phases.

For example, to get the  $\Delta_f G^{\text{Fe}_5\text{Th}}$ , the regression method was used with following equations.

$$\Delta_f G^{\text{Fe}_5\text{Th}}(T) = -19205 + 8.835T \quad (928 \text{ K} < T < 1164 \text{ K}) \quad (25)$$

$$\Delta_f G^{\text{Fe}_5\text{Th}}(1485 \text{ K}) = -8201 \quad (26)$$

$$\Delta_f G^{\text{Fe}_5\text{Th}}(1473 \text{ K}) = -8636 \quad (27)$$

where unit is J/mol.

As a result,  $\Delta_f G^{\text{Fe}_5\text{Th}}(T) = -14996 + 4.57T$  was

derived. The Gibbs energies of formation for the intermediate phases that were derived by the regression method are listed in Table 9. Final parameters defining the Gibbs energies of formation for the various phases are given in Table 10 and the phase diagram that was calculated from these parameters is shown in Figure 4. Comparison of calculated and experimental diagrams is done by the superposition in Figure 5.

The interaction parameters of the liquid phase combined to approach -30 kJ/mol for the heat of solution of Fe in Th at infinite dilution and -34 kJ/mol for the heat of solution of Th in Fe at infinite dilution. Refined heats of formation of compound phase are compared with those from experiment and from the Miedema approximation in Table 11.

### SUMMARY AND CONCLUSIONS

It was not the purpose of this study to reevaluate the phase diagram, but rather to generate internally consistent thermodynamics functions for each phase. From the phase diagram information, mostly invariant reactions, experimental data and a theoretical model, such thermodynamic descriptions for the phases were obtained with input information. First, all available thermodynamic data were assembled and a model which satisfactorily fits the experimental data was developed. Since any regression creates a best fit to all data, parameters were reevaluated such that they predicted certain critical phase diagram areas exactly along with representing the thermodynamic property. The second and third step in this study was set up in a way that each step optimized the parameters such that the difference between evaluated phase diagram and calculated one is approximately zero.

This analysis indicates that thermodynamic data alone can not be relied upon to predict the phase boundary behavior. Both thermodynamic data and phase diagram data should be used.

Table 1. Crystal Structure in the Fe-Th System

Phase	Proto type	Pearson Symbol	Space Group
$\alpha$ Th	Cu	cF4	Fm $\bar{3}$ m
$\gamma$ Fe	Cu	cF4	Fm $\bar{3}$ m
$\beta$ Th	W	cI2	Im $\bar{3}$ m
$\delta$ Fe	W	cI2	Im $\bar{3}$ m
$\alpha$ Fe	W	cI2	Im $\bar{3}$ m
Fe <sub>17</sub> Th <sub>2</sub>	Fe <sub>17</sub> Th <sub>2</sub>	hR19	R $\bar{3}$ m
Fe <sub>5</sub> Th	CaCu <sub>5</sub>	hP6	P6/mmm
$\alpha$ Fe <sub>7</sub> Th <sub>2</sub>	Ce <sub>2</sub> Ni <sub>7</sub>	hP36	P6 <sub>3</sub> /mmc
$\beta$ Fe <sub>7</sub> Th <sub>2</sub>	Gd <sub>2</sub> Co <sub>7</sub>	(rhombohedral)	—————
Fe <sub>3</sub> Th	NbBe <sub>3</sub>	hR12	R $\bar{3}$ m
Fe <sub>3</sub> Th <sub>7</sub>	Fe <sub>3</sub> Th <sub>7</sub>	hP20	P6 <sub>3</sub> mc

Table 2a. Lattice Stability of Fe and Th

Element	$\Delta_M G$ (J/mol)	Transformation Temperature (°C)
Th ( $\alpha$ - $\beta$ )	3600-2.2T	1360
( $\beta$ -L)	13800-6.8T	1735
( $\alpha$ -L)	17400-9.0T	1660
Fe ( $\alpha$ - $\gamma$ )	900-0.76T	911
( $\gamma$ - $\delta$ )	840-0.50T	1392
( $\delta$ -L)	13800-7.63T	1536
( $\gamma$ -L)	14640-8.13T	1528
( $\alpha$ -L)	15540-8.89T	1475

Table 2b. Gibbs Formation Energies of Intermediate Phases<sup>a</sup>

Phase	$\Delta_f G$ (J/mol)	Temperature Range (°C)
Fe <sub>17</sub> Th <sub>2</sub>	-12468+4.98T	(655 - 891 )
Fe <sub>5</sub> Th	-19205+8.835T	(655 - 891 )
Fe <sub>7</sub> Th <sub>2</sub>	-22845+11.21T	(655 - 891 )
Fe <sub>3</sub> Th	-24769+12.55T	(655 - 891 )
Fe <sub>3</sub> Th <sub>7</sub>	-5021-1.67T	(655 - 891 )

<sup>a</sup>Compound reaction  $x\text{Fe}(\alpha) + (1-x)\text{Th}(\alpha) \rightarrow \text{Fe}_x\text{Th}_{1-x}$ .

Table 3. Special Reactions of the Fe-Th System

Reaction	Compositions (at.% Th )			Temp. (°C)	Reaction Type
$\text{Fe}_{17}\text{Th}_2(\text{L}) \rightarrow \text{Fe}_{17}\text{Th}_2$	10.53			1462	Congruent
$\text{L} \rightarrow \delta\text{Fe} + \text{Fe}_{17}\text{Th}_2$	6.0	0	10.53	1412	Eutectic
$\text{L} + \text{Fe}_{17}\text{Th}_2 \rightarrow \text{Fe}_5\text{Th}$	26	10.53	16.7	1212	Peritectic
$\text{L} + \text{Fe}_5\text{Th} \rightarrow \text{Fe}_3\text{Th}$	30	16.7	25	1200	Peritectic
$\text{L} \rightarrow \text{Fe}_3\text{Th} + \text{Fe}_3\text{Th}_7$	55	25	70	875	Eutectic
$\text{L} + \alpha\text{Th} \rightarrow \text{Fe}_3\text{Th}_7$	67.5	99	70	962	Peritectic
$\text{Fe}_5\text{Th} + \text{Fe}_3\text{Th} \rightarrow \text{Fe}_7\text{Th}_2$	16.7	25	22.22	----	Peritectoid

Table 4. Enthalpy Differences between hcp, bcc and fcc Phases<sup>a</sup>

Z	3	4	5	5.5	6	7	8	8.5	9	10
hcp	-2.5	-2.5	10	15	13	-5	-10.5	-11	-8	-1
fcc	-2	-1.5	9	14	11	-3	-9.5	-11	-9	-2
bcc	2.2	2	-9.5	-14.5	-12	4	10	11	8.5	1.5

<sup>a</sup>Unit: KJ/mol.

Table 5. Change of the Structure Dependent Energy with the Average Number of Valence Electrons per Atom (i.e.,  $dE/dZ$ )<sup>a</sup>

Z		3	4	5	6	7	8	9	10
hcp	Z<0	-	0	13	-4	-18	-6	6	7
	Z>0	0	13	10	-18	-6	-1	7	-
fcc	Z<0	-	1	11	-6	-14	-7	4	7
	Z>0	1	11	10	-14	-7	-3	7	-
bcc	Z<0	-	0	-12	5	16	6	-5	-7
	Z>0	0	-12	-10	16	-6	2	-7	-

<sup>a</sup>Unit: KJ/mol per electron.



Table 6. Gibbs Energies Determined from Step 1 (Miedema model)<sup>a</sup>

Phase	$\Delta_M G_{Th}$	$\Delta_M G_{Fe}$	A	B
bcc1( $\alpha$ )	3599-2.204T	0	167000	-67000
bcc2( $\delta$ )	3599-2.204T	1740-1.2646T	167000	-67000
fcc	0	900-0.7061T	235000	-107000
liquid	17406-9.029T	15540-8.8976T	-42739	11756
Fe <sub>3</sub> Th <sub>7</sub>	-	-	-5021	1.67
Fe <sub>3</sub> Th	-	-	-24769	-12.55
Fe <sub>7</sub> Th <sub>2</sub>	-	-	-22845	-11.21
Fe <sub>5</sub> Th	-	-	-19205	-8.835
Fe <sub>17</sub> Th <sub>2</sub>	-	-	-12468	-4.98

<sup>a</sup>For solution phases:  $\Delta_M G^{xs} = X_{Th}(1-X_{Th})(A+B(2X_{Th}-1))$ ;  
 for solid intermediate phases:  $\Delta_f G = A - BT$ .

Table 7. Thermodynamic Information from Invariant Reactions<sup>a</sup>

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$\Delta_f G^{\text{Fe}_5\text{Th}}$ (1485) = $(d\Delta_M G^L/dx_{\text{Th}})$ (1485, 0.26)	
$(0.1667 - 0.26) + \Delta_M G^L(1485, 0.26)$	(1)
$\Delta_f G^{\text{Fe}_5\text{Th}}$ (1473) = $(d\Delta_M G^L/dx_{\text{Th}})$ (1473, 0.3 )	
$(0.1667 - 0.3 ) + \Delta_M G^L(1473, 0.3 )$	(2)
$\Delta_f G^{\text{Fe}_3\text{Th}}$ (1473) = $(d\Delta_M G^L/dx_{\text{Th}})$ (1473, 0.3 )	
$(0.25 - 0.3 ) + \Delta_M G^L(1473, 0.3 )$	(3)
$\Delta_f G^{\text{Fe}_3\text{Th}}$ (1148) = $(d\Delta_M G^L/dx_{\text{Th}})$ (1148, 0.55)	
$(0.25 - 0.55) + \Delta_M G^L(1148, 0.55)$	(4)
$\Delta_f G^{\text{Fe}_3\text{Th}_7}$ (1148) = $(d\Delta_M G^L/dx_{\text{Th}})$ (1148, 0.55)	
$(0.7 - 0.55) + \Delta_M G^L(1148, 0.55)$	(5)
$\Delta_f G^{\text{Fe}_3\text{Th}_7}$ (1235) = $(d\Delta_M G^L/dx_{\text{Th}})$ (1235, 0.675)	
$(0.7 - 0.675) + \Delta_M G^L(1235, 0.675)$	(6)
$\Delta_M G^{\text{fcc}}$ (1235, 0.99) = $(d\Delta_M G^L/dx_{\text{Th}})$ (1235, 0.675)	
$(0.7 - 0.675) + \Delta_M G^L(1235, 0.675)$	(7)
$\Delta_M G^{\text{bcc}}$ (1685, 0.0) = $(d\Delta_M G^L/dx_{\text{Th}})$ (1685, 0.1053)	
$(0 - 0.1053) + \Delta_M G^L(1685, 0.1053)$	(8)
$\Delta_f G^{\text{Fe}_{17}\text{Th}_2}$ (1735) = $(d\Delta_M G^L/dx_{\text{Th}})$ (1735, 0.1053)	
$(0) + \Delta_M G^L(1735, 0.1053)$	(9)

---

<sup>a</sup>G( T( K), mole fraction of Th).

Table 7. (Continued)

$$\Delta_f G^{\text{Fe}_{17}\text{Th}_2}(1685) = (d\Delta_M G^L/dx_{\text{Th}})(1685, 0.06) \\ (0.1053 - 0.06) + \Delta_M G^L(1685, 0.06) \quad (10)$$

$$\Delta_f G^{\text{Fe}_{17}\text{Th}_2}(1485) = (d\Delta_M G^L/dx_{\text{Th}})(1485, 0.26) \\ (0.1053 - 0.26) + \Delta_M G^L(1485, 0.26) \quad (11)$$


---

Table 8. Gibbs Energies of Formation Determined by  
Invariant Reaction (Step 2)<sup>a</sup>

Phase	$\Delta_M G_{Th}$	$\Delta_M G_{Fe}$	A	B
bcc1( $\alpha$ )	3599-2.204T	0	30000	0
bcc2( $\delta$ )	3599-2.204T	1740-1.2646T	30000	0
fcc	0	900-0.7061T	30000	0
liquid	17406-9.029T	15540-8.8976T	-32000	2000
Fe <sub>3</sub> Th <sub>7</sub>	-	-	-5021	1.67
Fe <sub>3</sub> Th	-	-	-24769	-12.55
Fe <sub>7</sub> Th <sub>2</sub>	-	-	-22845	-11.21
Fe <sub>5</sub> Th	-	-	-19205	-8.835
Fe <sub>17</sub> Th <sub>2</sub>	-	-	-12468	-2.274

<sup>a</sup>For solution phases:  $\Delta_M G^{xs} = X_{Th}(1-X_{Th})(A+B(2X_{Th}-1))$ ;  
for solid intermediate phases:  $\Delta_f G = A - BT$ .

Table 9. Gibbs Energies of Formation used in Last Step

Phase	$\Delta_f G$ (J/mol)	Temp. Range ( K)
$\text{Fe}_5\text{Th}$	$-19205 + 8.835T$	$928 < T < 1164$
	-8021	1485
	-8636	1473
$\text{Fe}_{17}\text{Th}_2$	$-12468 + 4.98T$	$928 < T < 1164$
	-7743	1735
	-7444	1685
	-6587	1485
$\text{Fe}_3\text{Th}_7$	$-5021 - 1.67T$	$928 < T < 1164$
	-7113	1235
	-6999	1148
$\text{Fe}_3\text{Th}$	$-24769 + 12.55T$	$928 < T < 1164$
	-10384	1473
	-10509	1148

Table 10. Gibbs Energies Determined from Last Step<sup>a</sup>

Phase	$\Delta_M^G \text{Th}$	$\Delta_M^G \text{Fe}$	A	B
bcc1( $\alpha$ )	3599-2.204T	0	30000	0
bcc2( $\delta$ )	3599-2.204T	1740-1.2646T	30000	0
fcc	0	900-0.7061T	30000	0
liquid	17406-9.029T	15540-8.8976T	-32000	2000
Fe <sub>3</sub> Th <sub>7</sub>	-	-	-4961	1.74
Fe <sub>3</sub> Th	-	-	-17311	-5.33
Fe <sub>7</sub> Th <sub>2</sub>	-	-	-22845	-11.21
Fe <sub>5</sub> Th	-	-	-14996	-4.57
Fe <sub>17</sub> Th <sub>2</sub>	-	-	-7953	-0.181

<sup>a</sup>For solution phases:  $\Delta_M^{G^{xs}} = X_{\text{Th}}(1-X_{\text{Th}})(A+B(2X_{\text{Th}}-1))$ ;  
 for solid intermediate phases:  $\Delta_f G = A - BT$ .

Table 11. Comparison of the Heats of Formation for  
Compound Phases

Phase	Experimental	Miedema Model	Last step (Regressed)
$\text{Fe}_{17}\text{Th}_2$	-12468	-5824	-7953
$\text{Fe}_5\text{Th}$	-19205	-9177	-14996
$\text{Fe}_7\text{Th}_2$	-22845	-11852	-22845
$\text{Fe}_3\text{Th}$	-24769	-12978	-17311
$\text{Fe}_3\text{Th}_7$	-5021	-9461	-4961

Table 12. Values for the Molar Volume, the Density of Electrons, the Chemical Potential for Electron and acf

	$V(\text{cm}^3/\text{mol})$	$N_{\text{ws}}^{1/3} (\text{du}^{1/3})^*$	$\phi^* (\text{V})$	acf
Fe	7.094	1.77	4.93	0.04
Th	19.79	1.28	3.30	0.04

$$P = 0.147 \text{ eV} / (\text{V}^2 \text{cm}^2 \text{du}^{1/3})$$

$$Q = 9.4 \text{ eV} / (\text{du} \text{cm}^2)$$

$$R = 0 \text{ eV} / (\text{du}^{1/3} \text{cm}^2)$$

$$\text{where } 1\text{eV} = 1.602 * 10^{-19} \text{ J}$$

$$\text{du}(\text{density unit}) = 100 \text{ kg}^{1/2} \text{cm}^{-5/2}$$

$$\text{du is defined as } (\text{K/V})^{1/2}.$$

Table 13. Parameters used in equation (7)

	electronic structure	Z	V ( $\text{cm}^3/\text{mol}$ )	$\mu * 10^{-6}$ ( $\text{kg}/\text{cm}^2$ )	$K * 10^{-6}$ ( $\text{kg}/\text{cm}^2$ )
Fe	$3d^6 4s^2$	8	7.094	0.831	1.716
Th	$6d^2 7s^2$	4	19.79	0.284	0.5534



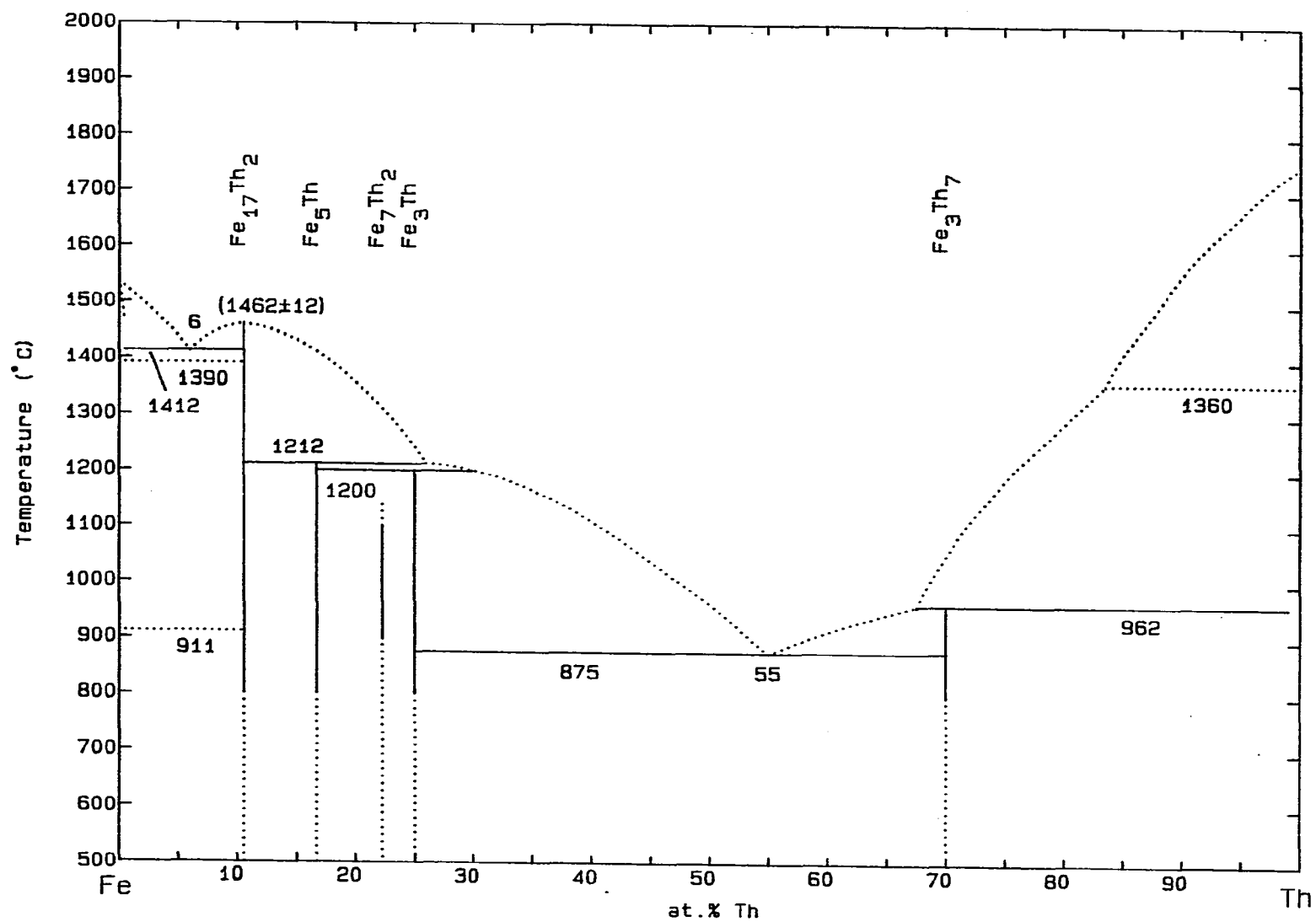


Fig. 1 The Fe-Th system evaluated from experiment

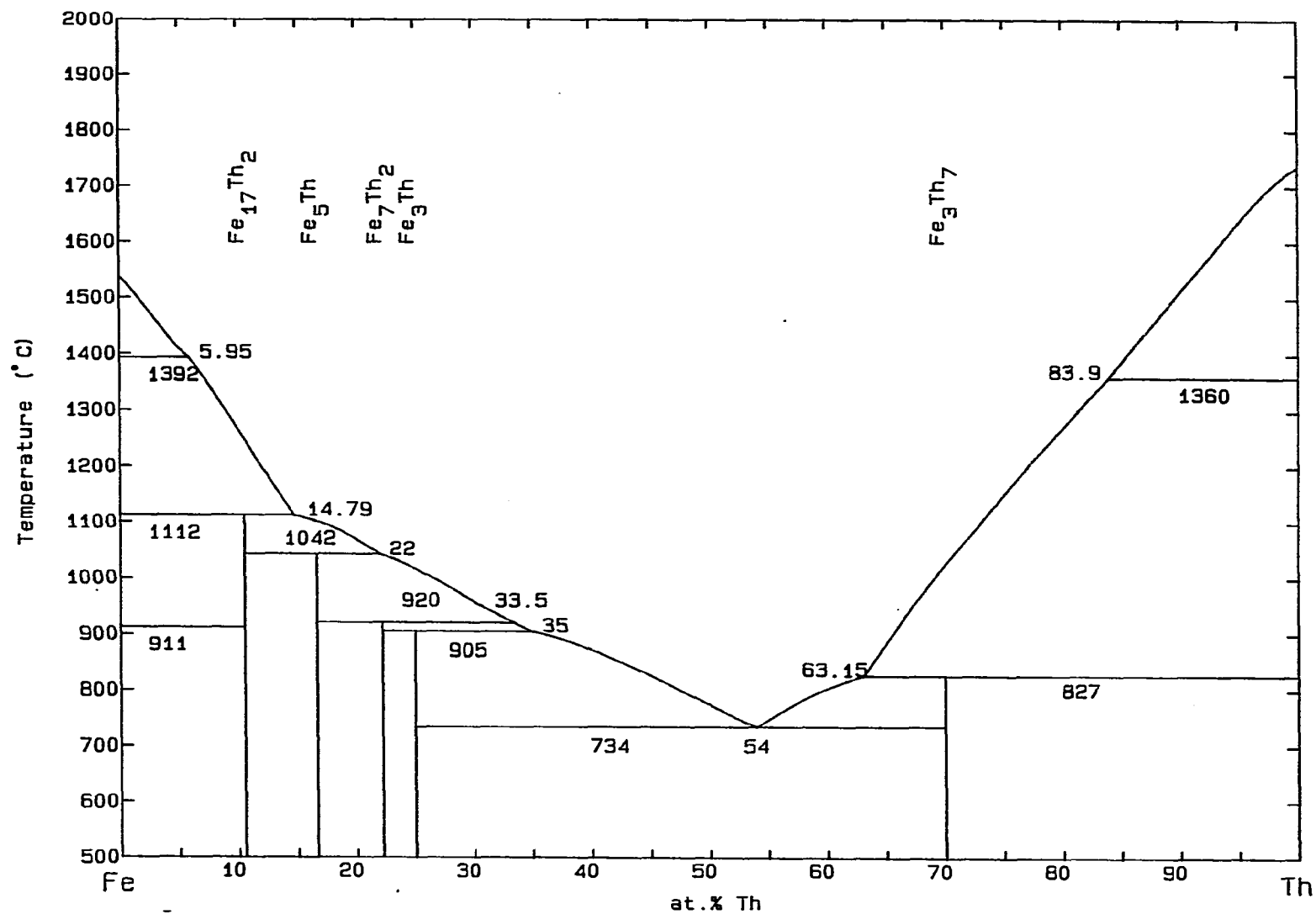


Fig. 2 The Fe-Th system (Miedema approximation)

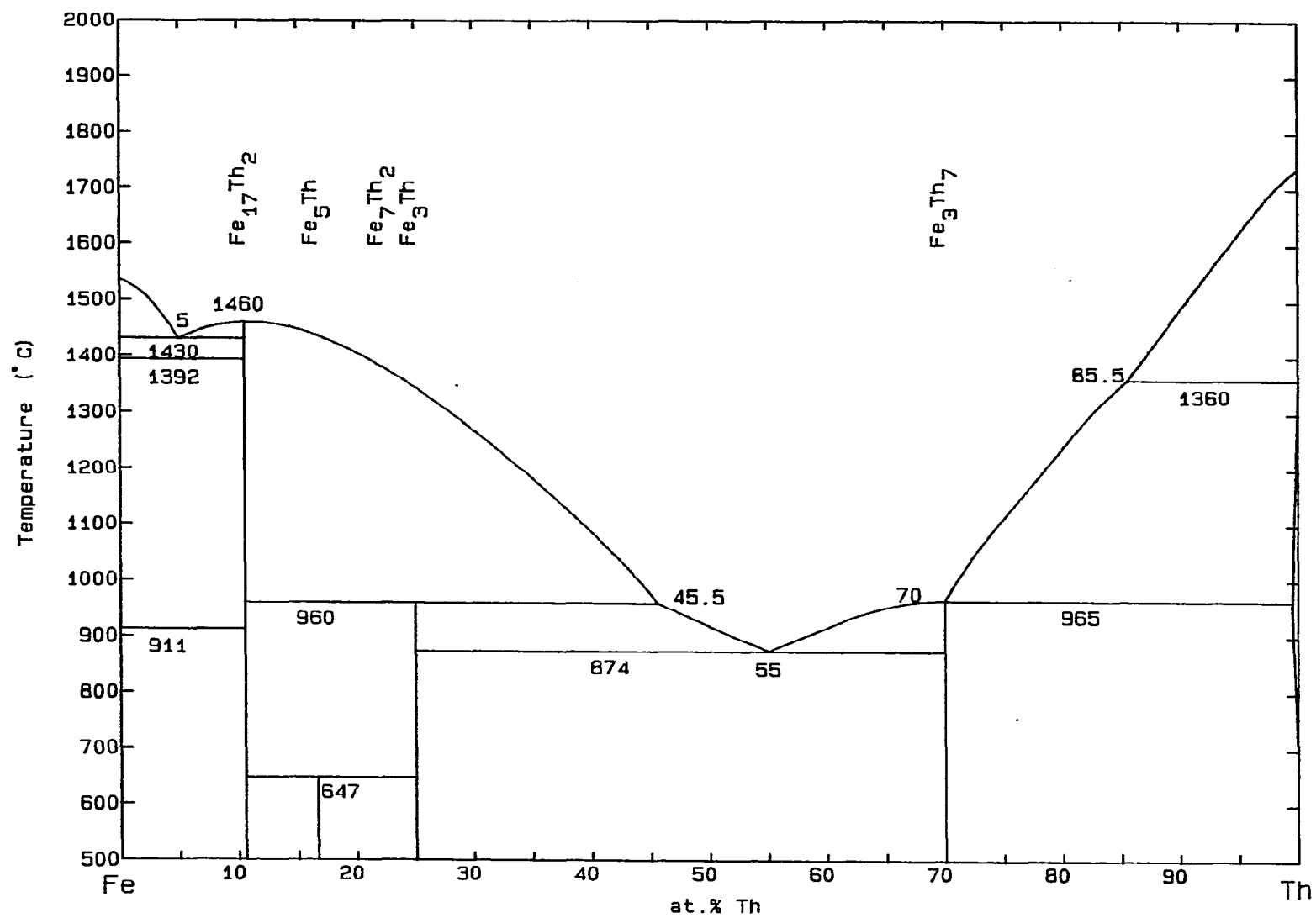


Fig. 3 The Fe-Th system (Second step analysis)

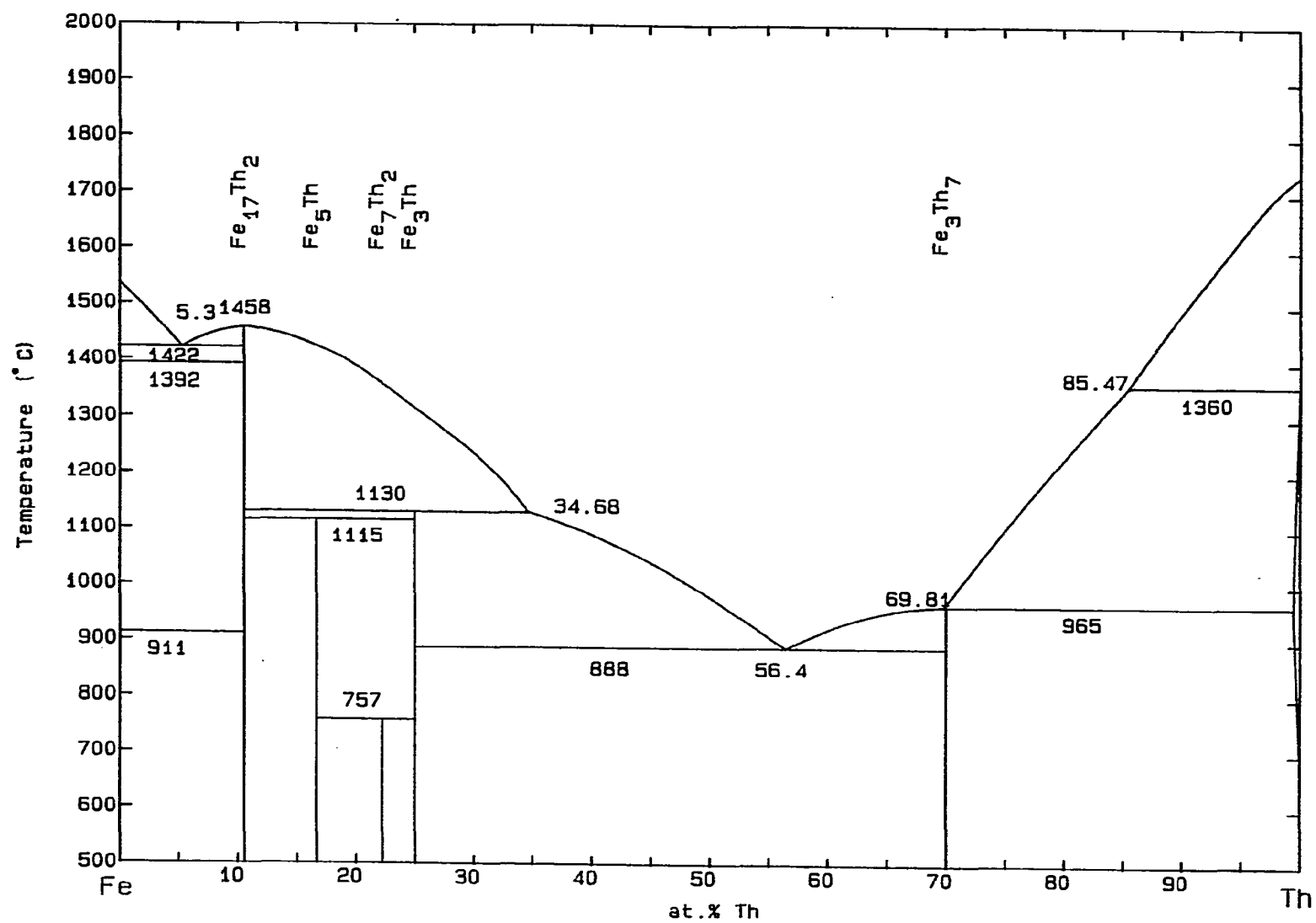


Fig. 4 The Fe-Th system (Last step analysis)

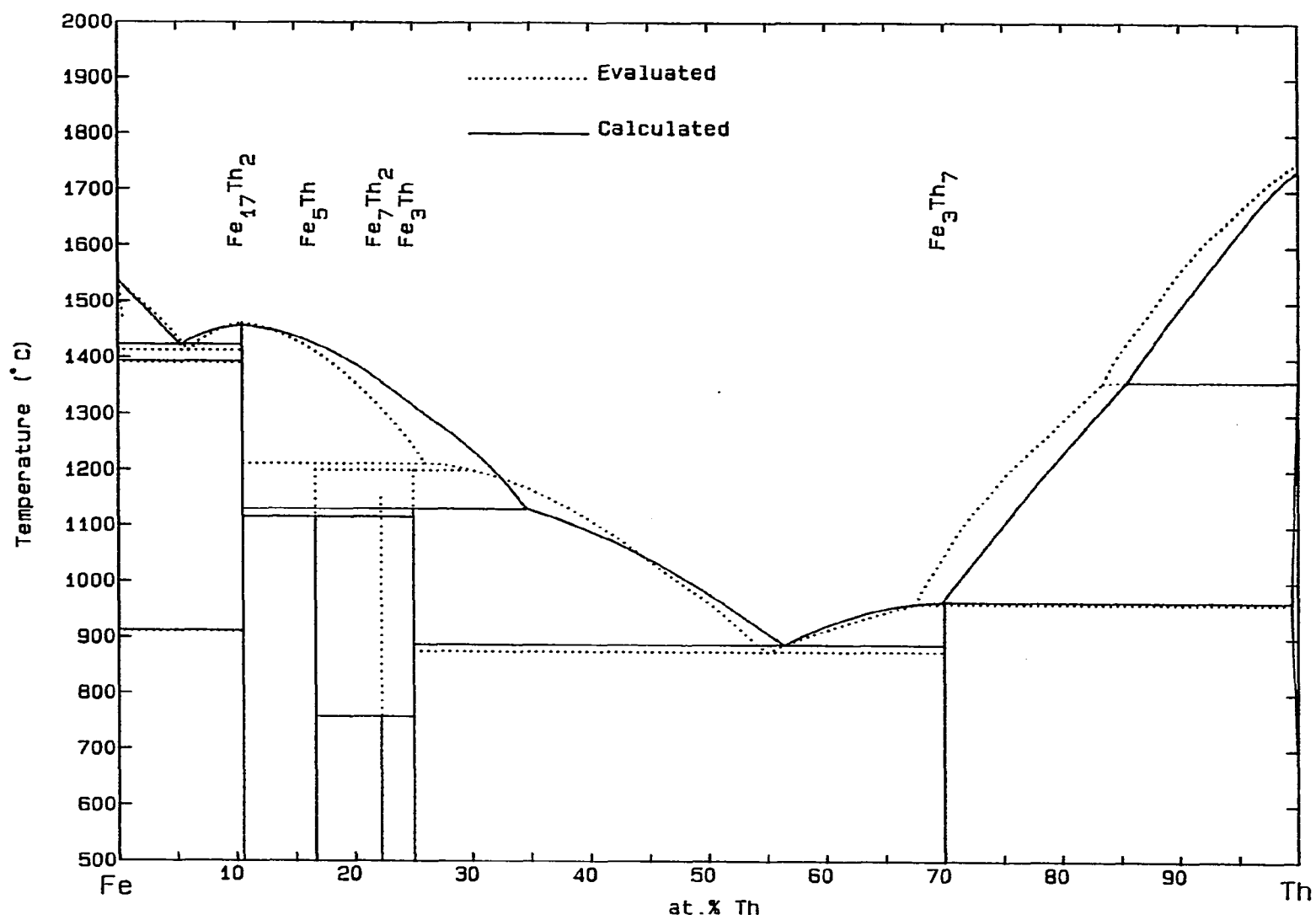


Fig. 5 The Fe-Th system (Comparison)

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PART II. GLASS FORMING ABILITY IN THE FE-TH SYSTEM

## INTRODUCTION

A glass is defined as a rigid body without long range atomic order. Glass formation is controlled by a combination of thermodynamic factors, kinetic factors and processing operations [Polk and Giessen 1978]. As a quantitative guide to the glass forming ability (GFA) in the Fe-Th system,  $T_0$  and TTT curves were evaluated from the previously derived thermodynamic description of the system.  $T_0$  curves, which are the loci of points defining  $\Delta_M G(\text{solid phase}) = \Delta_M G(\text{liquid phase})$ , are the thermodynamic limits for the partitionless crystallization of equilibrium phases. TTT curves, which result from a combined kinetic and thermodynamic approach, are time-temperature-transformation curves.

From the TTT curves, critical cooling rates ( $R_C$ ) to avoid crystallization were determined. Since the Fe-Th system has several solid phases, the relative stability of each phase in the metastable region is discussed. Kinetic parameters such as viscosity( $\eta$ ), glass transition temperature ( $T_g$ ) and interfacial energy ( $\sigma$ ) were estimated since experimental values were not available.



## THEORY

The determinations of  $T_0$  curves were carried out using the thermodynamic functions for the liquid phase and the terminal solid phases.  $T_0$  curves for intermediate phases were not determined because those phases were treated as invariant in composition.

To calculate the TTT curves, the Johnson-Mehl-Avrami treatment [Christian 1965] was used. According to their approach the crystallized volume fraction ( $x_f$ ) is

$$x_f = \frac{\pi}{3} I_v U_c^3 t^4 \quad (1)$$

where

$x_f$ : the transformed volume fraction

$I_v$ : the nucleation frequency

$U_c$ : the growth rate

$t$ : the time required to transform to solid

Nucleation was assumed to be homogeneous since glass transformation occurs when the cooling rate is high. The homogeneous nucleation frequency is

$$I_v = \frac{D_n N_v}{a_0} \exp \left( - \frac{\Delta_M G^*(T, X)}{kT} \right) \quad (2)$$

where

$D_n$  = the diffusion coefficient of liquid  
controlling nucleation

$N_v$  = the number of atoms per unit volume

$a_0$  = an atomic diameter

$\Delta_M G^*(T, X)$  = the Gibbs energy barrier for nucleation

By classical nucleation theory,  $\Delta_M G^*(T, X)$  is expressed

as

$$\Delta_M G^* = \frac{16\pi}{3N_a} \frac{\sigma_m^3}{G_m^2(T, X)} \quad (3)$$

where

$\sigma_m$ : the molar solid/liquid interfacial energy

$G_m$ : Gibbs energy of the liquid-crystal transformation

$N_a$ : Avogadro's number

Following Turnbull's approach [Turnbull 1950]  $\sigma_m = 0.43 \Delta_{fus} H$  was used where  $\Delta_{fus} H$  is the molar heat of fusion. The transformation energy was calculated with the following equation.

$$G_m(T, X) = G^L(T, X_L) - G^S(T, X_S) + \frac{dG^L(T, X_L)}{dX} (X_S - X_L) \quad (4)$$

where  $X$  is the mole fraction of Th. The terms in equation (4) are defined by illustration in the 700°C G-X diagram of the system (Figure 6) with solid phase being  $Fe_3Th_7$ ,

$T = 700^{\circ}\text{C}$  and  $X = X_L = 0.6$ .

The growth rate of a nucleus is

$$U_c = \frac{f D_g}{a_0} (1 - \exp(-G_m(T, X)/RT)) \quad (5)$$

where

$D_g$ : the diffusion coefficient of liquid controlling growth

$f$ : the fraction of atomic sites at the liquid/crystal interface (0 - 1)

On the basis of the Stokes - Einstein equation,  $D_g$  and  $D_n$  were assumed as

$$D_n = D_g = D = kT/3\eta a_0 \quad (6)$$

where  $\eta$  is the viscosity of undercooled liquid and  $D$  is the bulk diffusion coefficient of liquid. Since the experimental viscosity of the liquid phase was not available between the melting temperature and the glass transition temperature, the free volume model suggested by Ramachandrarao et al. [1977] was used. The viscosity with the model is

$$\eta = \eta_0 \exp(1/f_v) \quad (7)$$

$$f_v = a_1 \exp(-E_h/RT) \quad (8)$$

$$E_h = 57.865 T_g - 5857.6 \text{ (J/mol)} \quad (9)$$

where  $f_v$  is the free volume fraction,  $E_h$  is the energy of vacancy formation and  $\eta_0$  and  $a_1$  are constants. Equation (9) was derived empirically by Ramachandrarao et al.  $\eta_0$  and  $a_1$  were evaluated by the condition that

$$\eta = 10^{12} \text{ N/msec and } f_v = 0.03 \text{ when } T = T_g \quad (10)$$

which is considered to be a characteristic of glass transition materials. In equation (10)  $T_g$  is the glass transition temperature.

$T_g$  values as a function of composition were not available for the Fe-Th system. However for many glasses glass transition temperature vs. melting temperature ( $T_g/T_m$ ) which is called the reduced glass transition temperature is near constant and  $T_g$  is very similar to crystallization temperature ( $T_x$ ) [Reeve et al. 1982] though  $T_g$  is less sensitive to the cooling rate ( $\dot{T}$ ) than  $T_x$  is [Massalski et al. 1983]. It was reported that  $T_g/T_m$  was about 0.25 for pure metals and 0.45 to 0.65 for easy glass forming alloys [Davies 1976]. For the calculation  $T_g/T_m$  values were assumed to be 0.25 at  $X = 0$  and 1 increasing to

0.5 at  $X = 0.5$ , so that

$$\begin{aligned} T_g/T_m &= 0.25 + 0.5X && \text{when } 0 < X < 0.5 \\ &= 0.25 + 0.5(1-X) && \text{when } 0.5 < X < 1 \end{aligned} \quad (11)$$

$T_g$  and  $T_m$  values with composition are shown in Figure 8.

From equations (1), (2), (5) and (7),  $t$  is given as

$$t = \frac{9.3}{kT} \left( \frac{a_0^9 x_f}{f^3 N_v} \frac{\exp(\Delta_M G^*(T, X)/kT)}{(1 - \exp(G_m/RT))^3} \right)^{1/4} \quad (12)$$

With equation (12), TTT curves were calculated for all solid phases. The critical cooling rate ( $R_c$ ) with composition change was derived by

$$R_c = \frac{T_m - T_n}{t_n} \quad (13)$$

where  $T_n$  = the nose temperature

$t_n$  = the time at the nose

## RESULTS AND DISCUSSION

 $T_0$  curves determination

For the calculation of  $T_0$  curves, it was necessary to evaluate the Gibbs energy of undercooled liquid over the range from glass transition temperature ( $T_g$ ) to the melting temperature ( $T_m$ ) [Sommer 1985]. Because of the lack of heat capacity ( $C_p$ ) data,  $\Delta C_p^{l-g}$  was assumed to be zero. Then, with the previously evaluated Gibbs energy functions for the solution phases,  $T_0$  curves for fcc, bccl and bcc2 were calculated and are shown in Figure 7. The resulting  $T_0$  curves, which imply a wide glass forming region in the system, fall quite rapidly so that partitionless crystallization of the terminal solid solutions is indicated to be very limited [Boettinger 1982a, 1982b]. Near room temperature (27°C),  $T_0$  occurs at 20.5 at.% Th and 79.2 at.% Th for bccl phase, 18.4 at.% Th and 79.8 at.% Th for bcc2 and 20.9 at.% Th and 74.3 at.% Th for fcc phase.

TTT curves determination

For the calculation of TTT curves,  $N_v$ ,  $a_0$ ,  $X_f$  and  $f$ , which are, respectively,  $5E28$  atoms/m<sup>3</sup>,  $3.6E-10$  m,  $1E-6$  and  $0.1$ , were used in equation (12). Typical TTT curves for the solid phases are shown in Figures 9-19 from 0 to 100 at.% Th with an interval of 10 at.% Th. Values from the

TTT curves for  $T_n$  and  $t_n$  were used to calculate the  $R_c$ 's. Characteristic quantities associated with the TTT curves of all solid phases in the system are given in Tables 14 through 21. With data from Tables 14 - 21, critical cooling rates were calculated and are shown in Figure 20. It may be noted that the critical cooling rate for each phase is maximum at its own stoichiometric composition because the maximum driving force toward crystallization exists at stoichiometry.  $Fe_7Th_2$  and  $Fe_5Th$  phases are less stable than the other phases so their critical cooling rate curves are less than those of other phases. The critical cooling rate curves of the terminal solid solution phases (bcc1( $\alpha$ ), bcc2( $\delta$ ) and fcc) are quite similar to the  $T_0$  curves. Table 22 shows data for the phases in which glass formation would be most difficult. With a cooling rate of  $1E6$  (K/sec), the possible glass forming region is 45 - 52 at.% Th which is consistent with glass formation that has been observed at 50 at.% Th by the melt spinning process ( $\dot{T} = 1E6$  K/sec) [Drehman, private communication 1985]. The glass forming region determined from  $T_0$  curves is compared with that from TTT curves in Figure 21. Figure 21 shows that at high cooling rates ( $\dot{T} > 1E9$  K/sec), glass formation is competing with composition-invariant phase crystallization [Massalski et al. 1983]. However glass can be formed outside the glass forming regions that are

determined by  $T_0$  curves if the cooling rates are above  $1E10$  K/sec [Lin et al. 1984]. It appears to be possible to form amorphous alloys throughout the whole composition range in this system if the cooling rates exceed  $1E12$  K/sec.



## CONCLUSIONS

The glass forming ability of Fe-Th alloys was determined from TTT and  $T_0$  curves with data for the thermodynamic and kinetic parameters being based upon input from the calculation of the equilibrium diagram. It has thus been demonstrated that the Gibbs energy functions describing the equilibrium phase diagram can be useful in deriving metastable as well as stable phase formation, and a complete assessment of the metastable equilibria should include consideration of equilibrium thermodynamic functions as well as kinetic factors.

Table 14. Characteristics of TTT Curves for bccl Phase

at.% Th	$T_n$ ( K)	$t_n$ (sec)	$R_c$ ( K/sec)	$\log_{10}(R_c)$
0	965	0.39E-8	0.21E12	11.33
5	860	0.21E-7	0.41E11	10.60
10	770	0.19E-5	0.49E9	8.69
15	660	0.63	0.16E4	3.21
20 - 75	-	-	-	< -17
80	530	0.11E16	0.86E12	-12.06
85	690	0.31E-2	0.31E6	5.49
90	845	0.32E-6	0.30E10	9.48
95	990	0.82E-8	0.11E12	11.06
100	1125	0.13E-8	0.67E12	11.82

Table 15. Characteristics of TTT Curves for fcc Phase

at.% Th	$T_n$ ( K)	$t_n$ (sec)	$R_c$ ( K/sec)	$\log_{10}(R_c)$
0	990	0.26E-8	0.31E12	11.49
5	870	0.14E-7	0.59E11	10.76
10	780	0.12E-5	0.73E9	8.66
15	665	0.31	0.32E4	3.51
20 - 70	-	-	-	< -17
75	520	0.36E20	0.25E-16	-16.58
80	635	0.40E1	0.22E3	2.35
85	755	0.32E-4	0.28E8	7.45
90	870	0.24E-6	0.39E10	9.59
95	970	0.16E-7	0.58E11	10.77
100	1060	0.33E-8	0.28E12	11.45

Table 16. Characteristics of TTT Curves for bcc2 Phase

at.% Th	$T_n$ ( K)	$t_n$ (sec)	$R_c$ ( K/sec)	$\log_{10}(R_c)$
0	1000	0.22E-8	0.36E12	11.56
5	870	0.12E-7	0.64E11	10.81
10	760	0.22E-5	0.43E9	8.63
15	620	0.29E3	0.37E1	0.57
20 - 75	-	-	-	< -17
80	510	0.84E20	0.12E-16	-16.91
85	685	0.55E-2	0.18E6	5.25
90	845	0.31E-6	0.31E10	9.49
95	990	0.79E-8	0.11E12	11.07
100	1125	0.13E-8	0.67E12	11.82

Table 17. Characteristics of TTT Curves for  $\text{Fe}_3\text{Th}_7$  Phase

at.% Th	$T_n$ (K)	$t_n$ (sec)	$R_c$ (K/sec)	$\log_{10}(R_c)$
0 -30	-	-	-	<-17
35	635	0.13E11	0.59E7	-7.22
40	695	0.66E4	0.10	-0.98
45	750	0.15E1	0.35E3	2.55
50	785	0.10E-1	0.44E5	4.64
55	780	0.21E-4	0.17E8	7.23
60	795	0.48E-5	0.81E8	7.91
65	795	0.17E-5	0.25E9	8.39
70	805	0.18E-5	0.27E9	8.44
75	815	0.39E-5	0.16E9	8.21
80	795	0.72E-5	0.10E9	8.02
85	770	0.35E-4	0.25E8	7.40
90	730	0.75E-3	0.14E7	6.16
95	660	0.34	0.37E4	3.57
100	-	-	-	< -17

Table 18. Characteristics of TTT Curves for Fe<sub>3</sub>Th Phase

at.% Th	T <sub>n</sub> ( K)	t <sub>n</sub> (sec)	R <sub>C</sub> ( K/sec)	log <sub>10</sub> (R <sub>C</sub> )
0	475	0.19E14	0.66E-10	-10.17
5	730	0.12E-4	0.80E8	7.90
10	805	0.87E-5	0.10E9	8.02
15	850	0.85E-5	0.10E9	8.00
20	870	0.81E-5	0.92E8	7.96
25	870	0.66E-5	0.97E8	7.98
30	880	0.13E-4	0.42E8	7.63
35	885	0.41E-4	0.13E8	7.12
40	880	0.13E-3	0.37E7	6.57
45	860	0.45E-3	0.99E6	5.99
50	835	0.19E-2	0.20E6	5.31
55	765	0.46E-3	0.82E6	5.91
60	725	0.28E-2	0.16E6	5.20
65	685	0.30E-1	0.17E5	4.25
70	805	0.86E1	0.77E2	1.88
75	610	0.55E6	0.15E-2	-2.81
80	560	0.81E13	0.12E-9	-9.91
85 - 100	-	-	-	<-17

Table 19. Characteristics of TTT Curves for  $\text{Fe}_7\text{Th}_2$  Phase

at.% Th	$T_n$ (K)	$t_n$ (sec)	$R_c$ (K/sec)	$\log_{10}(R_c)$
0	515	0.29E8	0.44E-4	-4.35
5	715	0.96E-4	0.10E8	7.00
10	775	0.13E-3	0.71E7	6.85
15	810	0.20E-3	0.43E7	6.63
20	825	0.26E-3	0.29E7	6.47
25	820	0.25E-3	0.27E7	6.44
30	830	0.76E-3	0.84E6	5.92
35	835	0.36E-2	0.16E6	5.22
40	830	0.17E-1	0.32E5	4.51
45	810	0.80E-1	0.62E4	3.79
50	790	0.44	0.99E3	2.99
55	725	0.44E-1	0.95E4	3.98
60	695	0.30	0.16E4	3.21
65	655	0.31E1	0.18E3	2.26
70	630	0.11E4	0.60	-0.21
75	605	0.69E8	0.12E-4	-4.90
80	560	0.11E15	0.84E-11	-11.07
85 - 100	-	-	-	<-17

Table 20. Characteristics of TTT Curves for Fe<sub>5</sub>Th Phase

at.% Th	T <sub>n</sub> ( K)	t <sub>n</sub> (sec)	R <sub>C</sub> ( K/sec)	log <sub>10</sub> (R <sub>C</sub> )
0	570	0.13E2	0.89E2	1.95
5	810	0.52E-6	0.16E10	9.22
10	870	0.64E-6	0.13E10	9.12
15	900	0.95E-6	0.83E9	8.92
20	910	0.14E-5	0.50E9	8.70
25	900	0.18E-5	0.34E9	8.53
30	895	0.53E-5	0.10E9	8.03
35	890	0.25E-4	0.21E8	7.33
40	870	0.14E-3	0.34E7	6.53
45	840	0.11E-2	0.40E6	5.60
50	805	0.15E-1	0.27E5	4.44
55	725	0.62E-2	0.67E5	4.83
60	680	0.21	0.23E4	3.37
65	625	0.31E2	0.18E2	1.27
70	580	0.78E7	0.94E-4	-4.02
75 - 100	-	-	-	<-17



Table 21. Characteristics of TTT Curves for  $\text{Fe}_{17}\text{Th}_2$  Phase

at.% Th	$T_n$ (K)	$t_n$ (sec)	$R_c$ (K/sec)	$\log_{10}(R_c)$
0	665	0.89E-4	0.12E8	7.10
5	925	0.15E-7	0.50E11	10.70
10	980	0.20E-7	0.36E11	10.56
15	995	0.33E-7	0.20E11	10.31
20	990	0.60E-7	0.10E11	10.02
25	960	0.10E-6	0.51E10	9.71
30	940	0.40E-6	0.13E10	9.11
35	910	0.29E-5	0.18E9	8.25
40	870	0.40E-4	0.12E8	7.10
45	820	0.15E-2	0.32E6	5.51
50	760	0.36	0.12E4	3.11
55	660	0.71E1	0.27E3	2.44
60	590	0.18E6	0.31E-2	-2.49
65	515	0.63E16	0.11E-12	-12.95
70 - 100	-	-	-	<-17

Table 22. Critical Cooling Rate in the Th-Fe System

at.% Th	most competing phase	$\log_{10} (R_c)$
0	bcc2	11.56
5	bcc2	10.81
10	$\text{Fe}_{17}\text{Th}_2$	10.56
15	$\text{Fe}_{17}\text{Th}_2$	10.31
20	$\text{Fe}_{17}\text{Th}_2$	10.02
25	$\text{Fe}_{17}\text{Th}_2$	9.71
30	$\text{Fe}_{17}\text{Th}_2$	9.11
35	$\text{Fe}_{17}\text{Th}_2$	8.25
40	$\text{Fe}_{17}\text{Th}_2$	7.10
45	$\text{Fe}_3\text{Th}$	5.99
50	$\text{Fe}_3\text{Th}$	5.31
55	$\text{Fe}_3\text{Th}_7$	7.23
60	$\text{Fe}_3\text{Th}_7$	7.91
65	$\text{Fe}_3\text{Th}_7$	8.39
70	$\text{Fe}_3\text{Th}_7$	8.44
75	$\text{Fe}_3\text{Th}_7$	8.21
80	$\text{Fe}_3\text{Th}_7$	8.02
85	fcc	7.45
90	fcc	9.59
95	bcc2	11.07
100	bcc1/bcc2	11.82

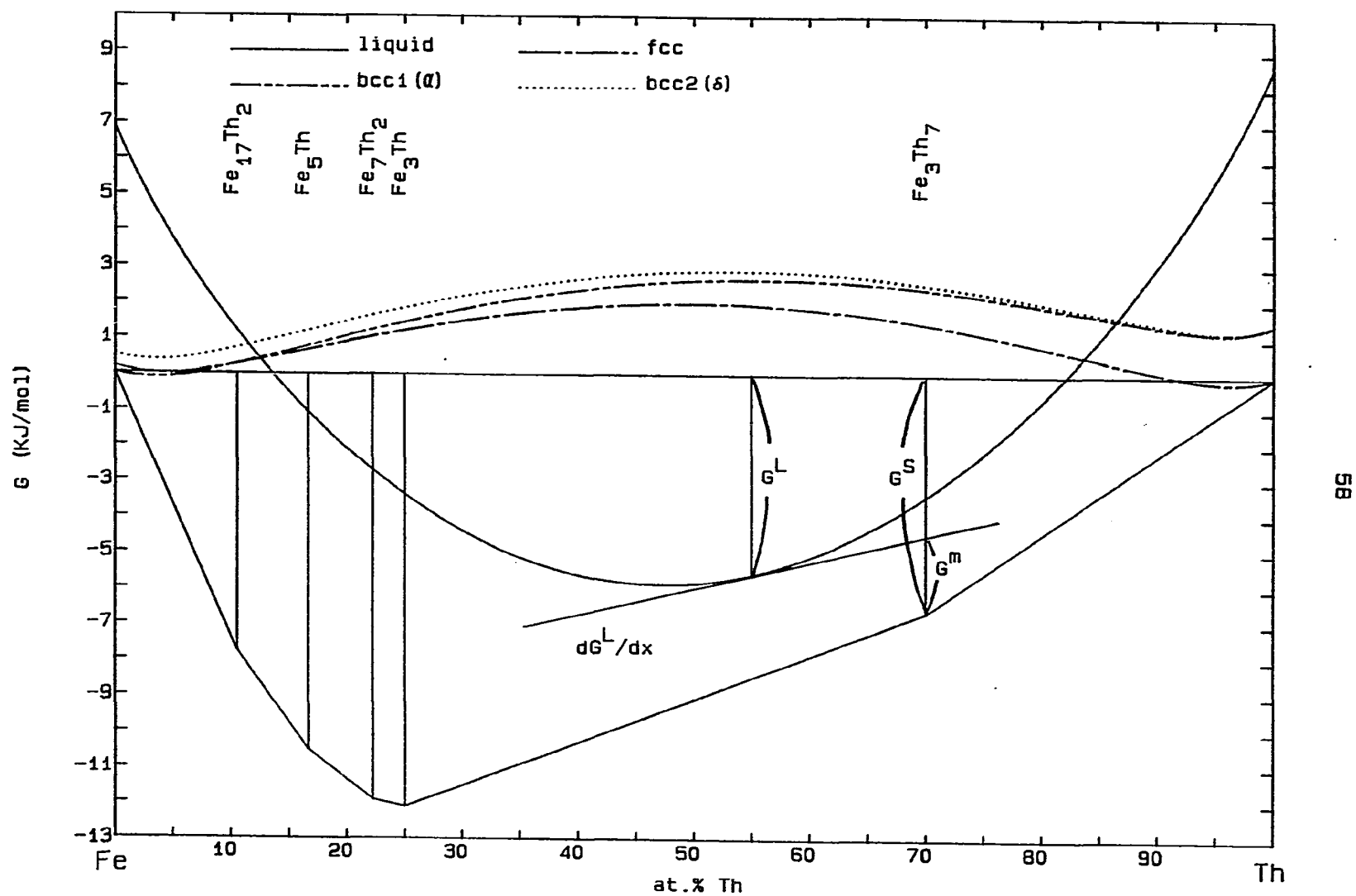


Fig. 6 The G-x diagram (700°C)

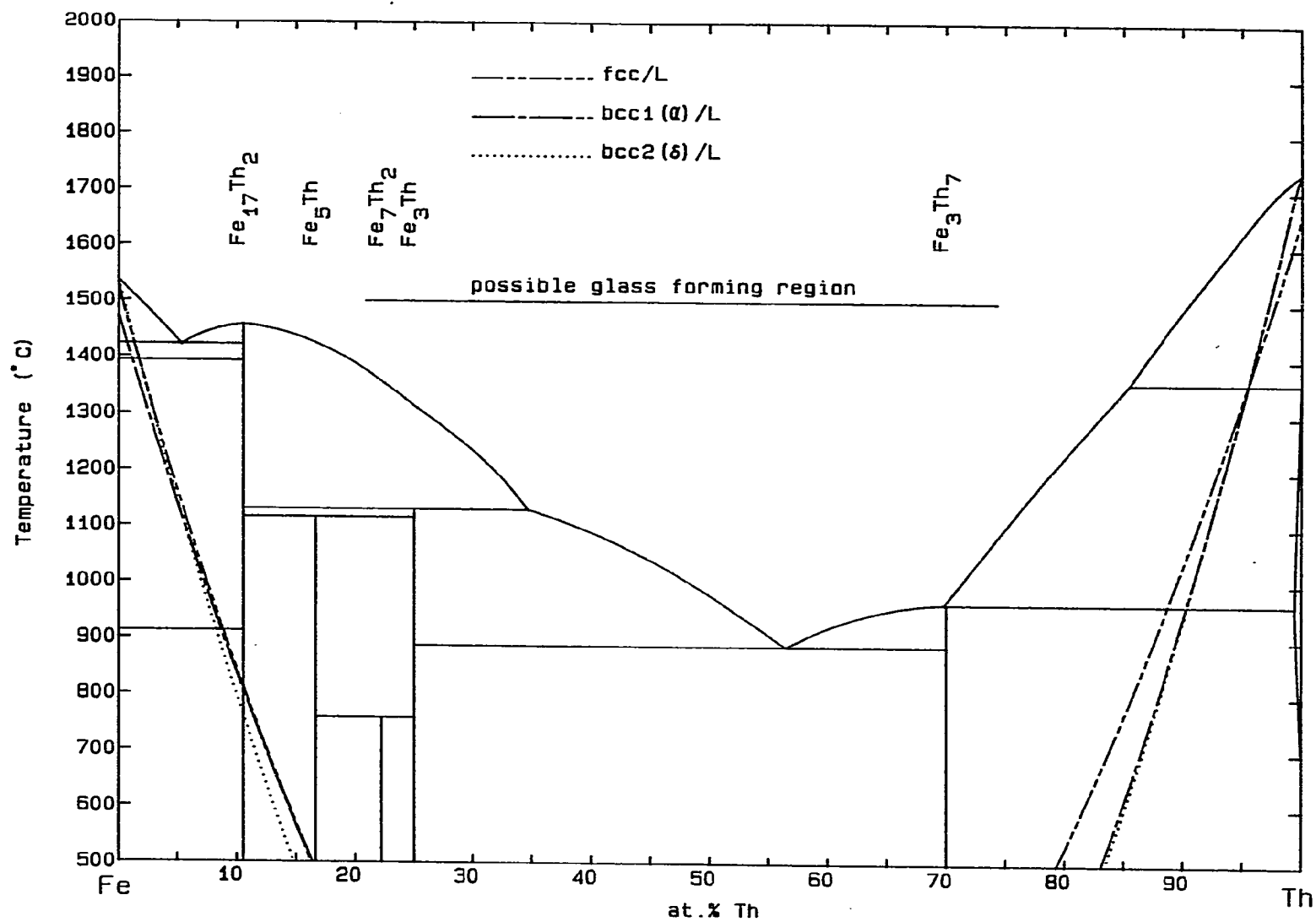


Fig. 7 The Fe-Th system ( $T_0$  Curves)

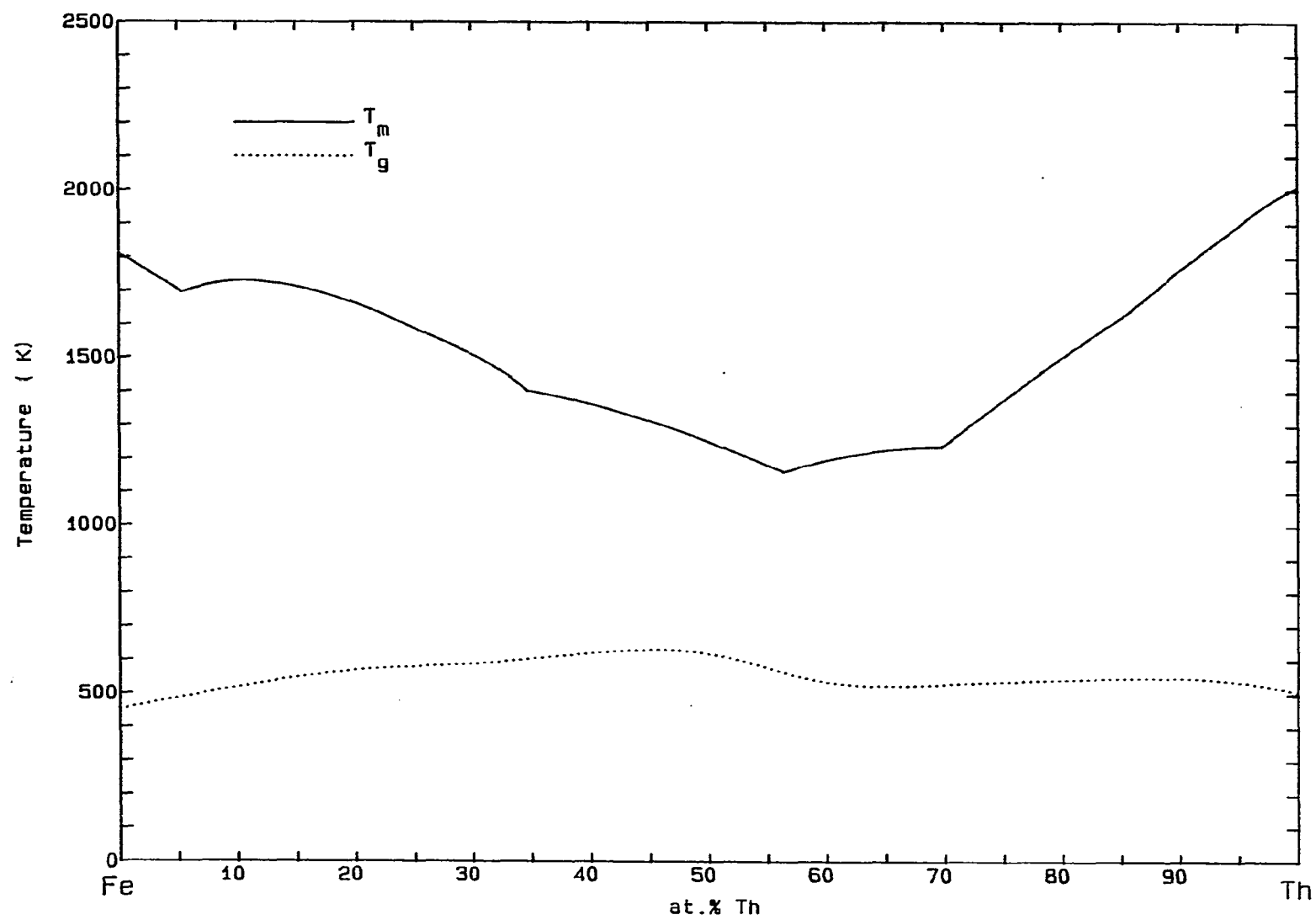


Fig. 8  $T_m$  and  $T_g$  in the Fe-Th system

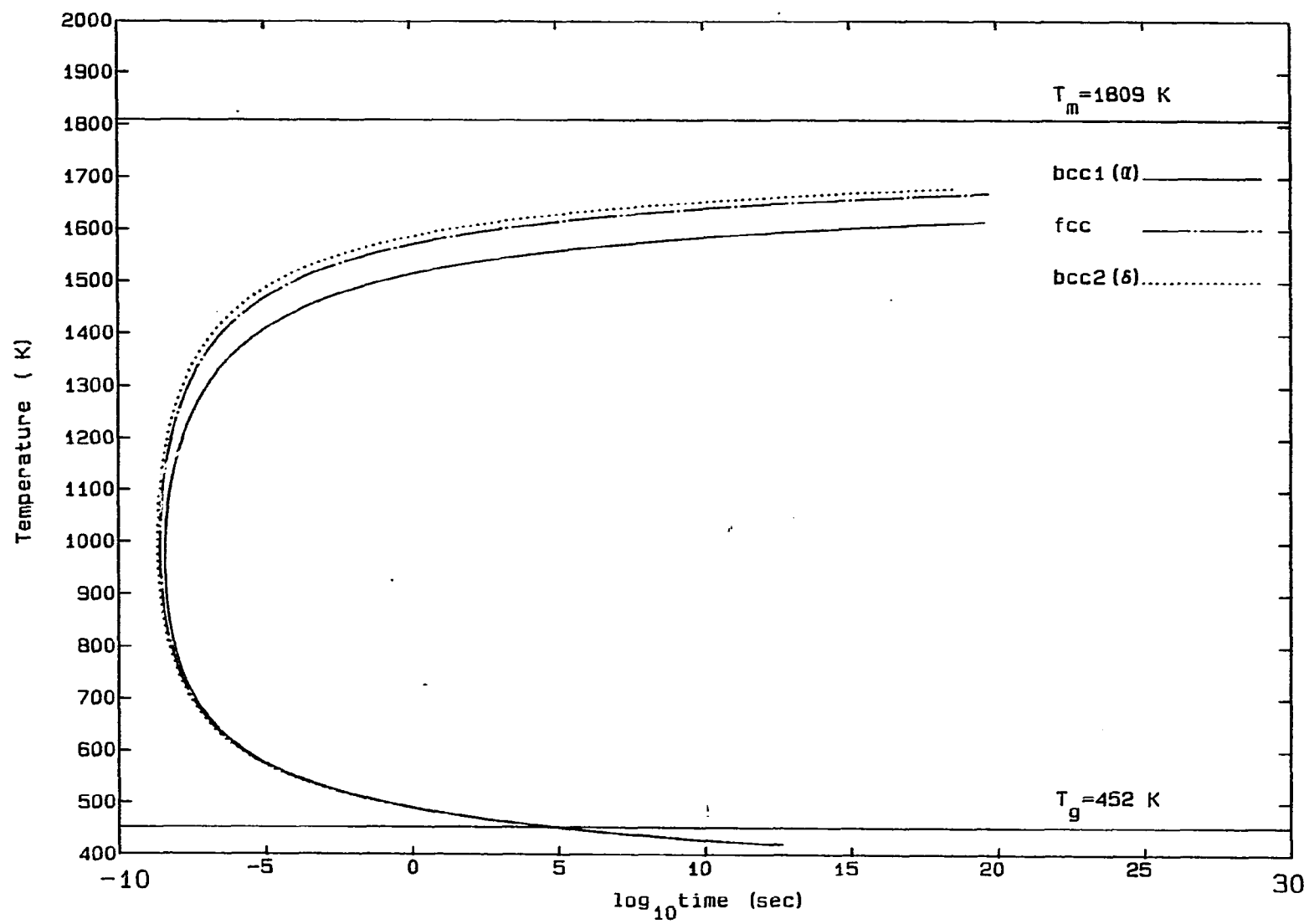


Fig. 9 TTT curves (0 at.% Th)

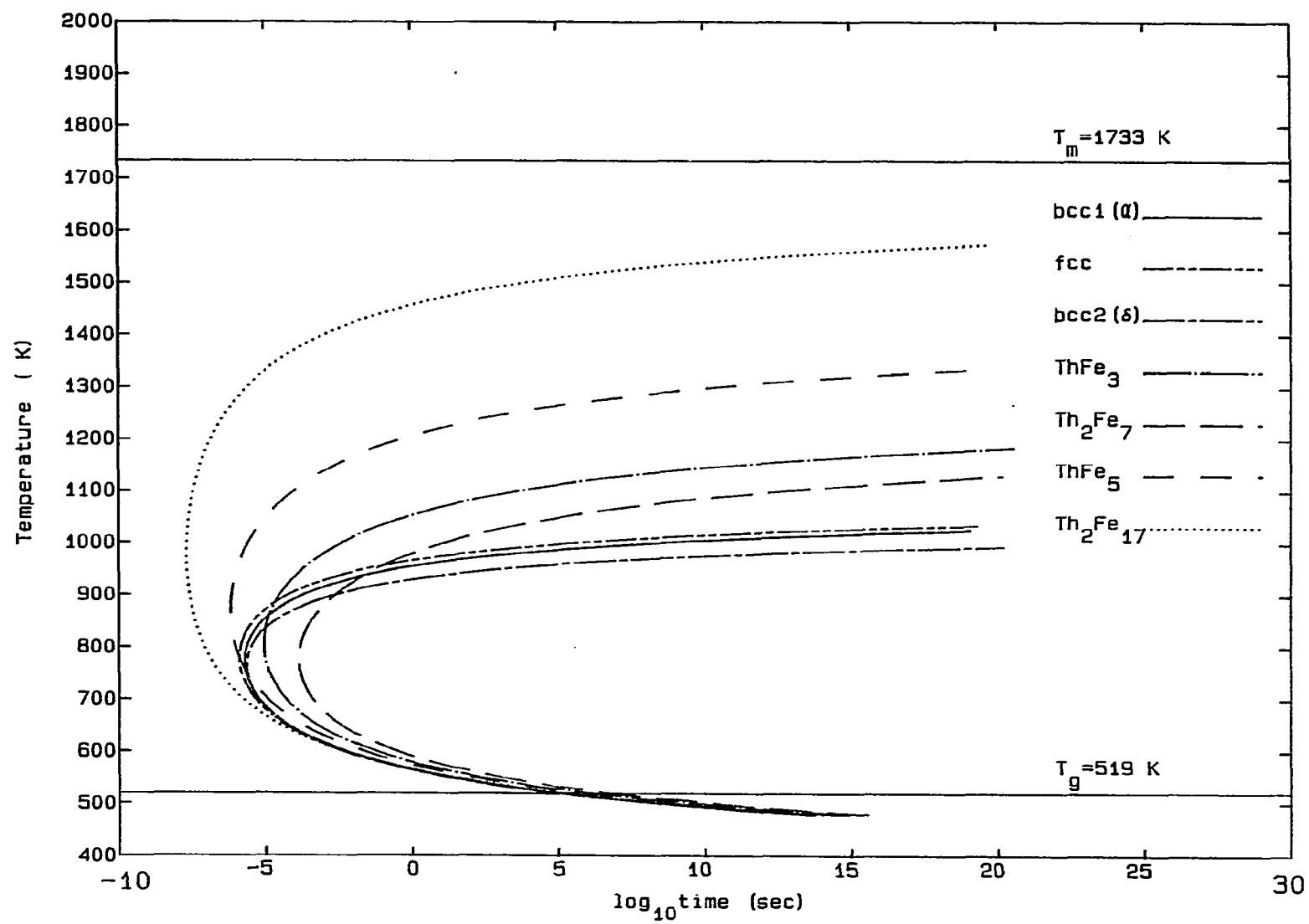


Fig. 10 TTT curves (10 at.% Th)

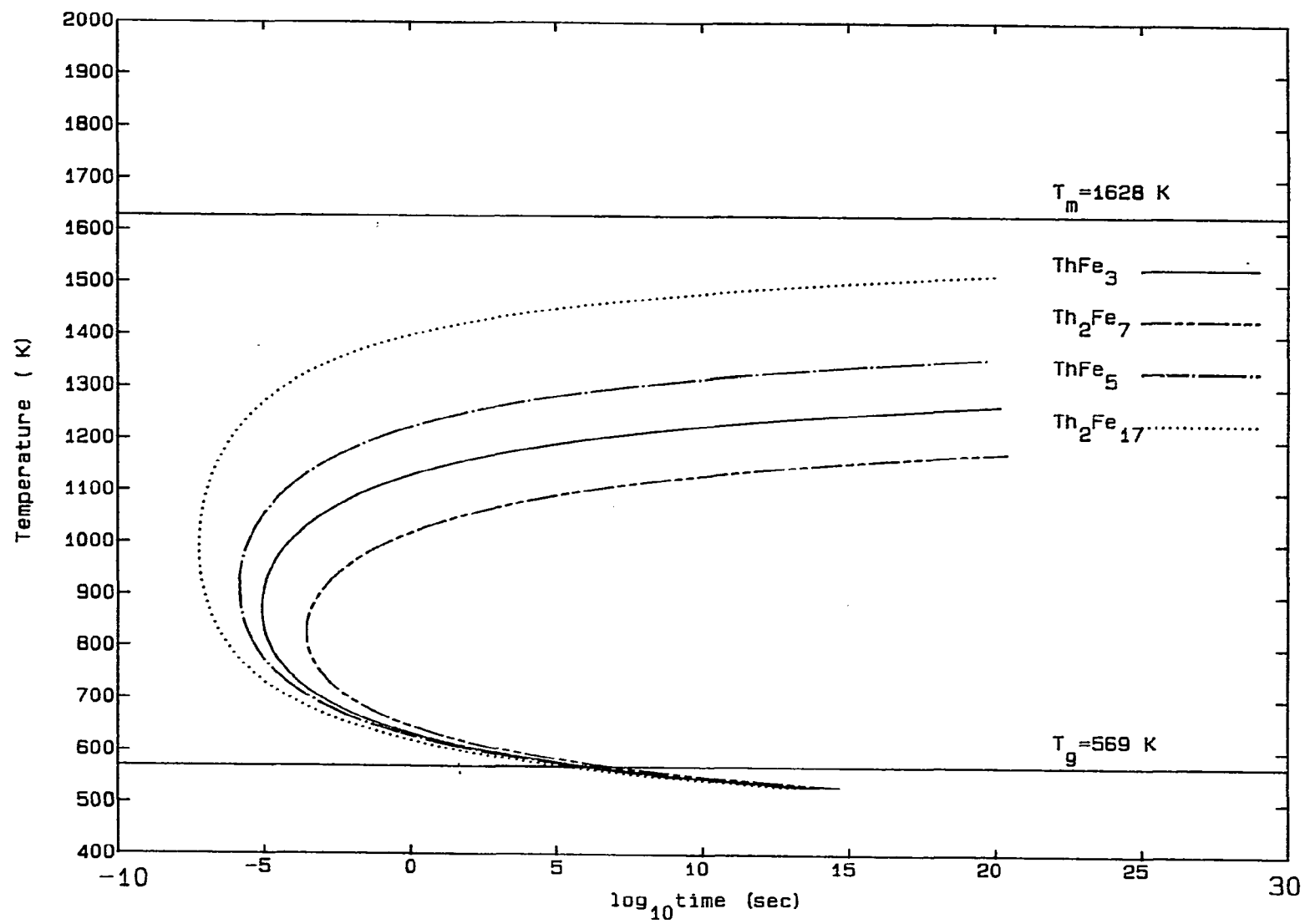


Fig. 11 TTT curves (20 at.% Th)



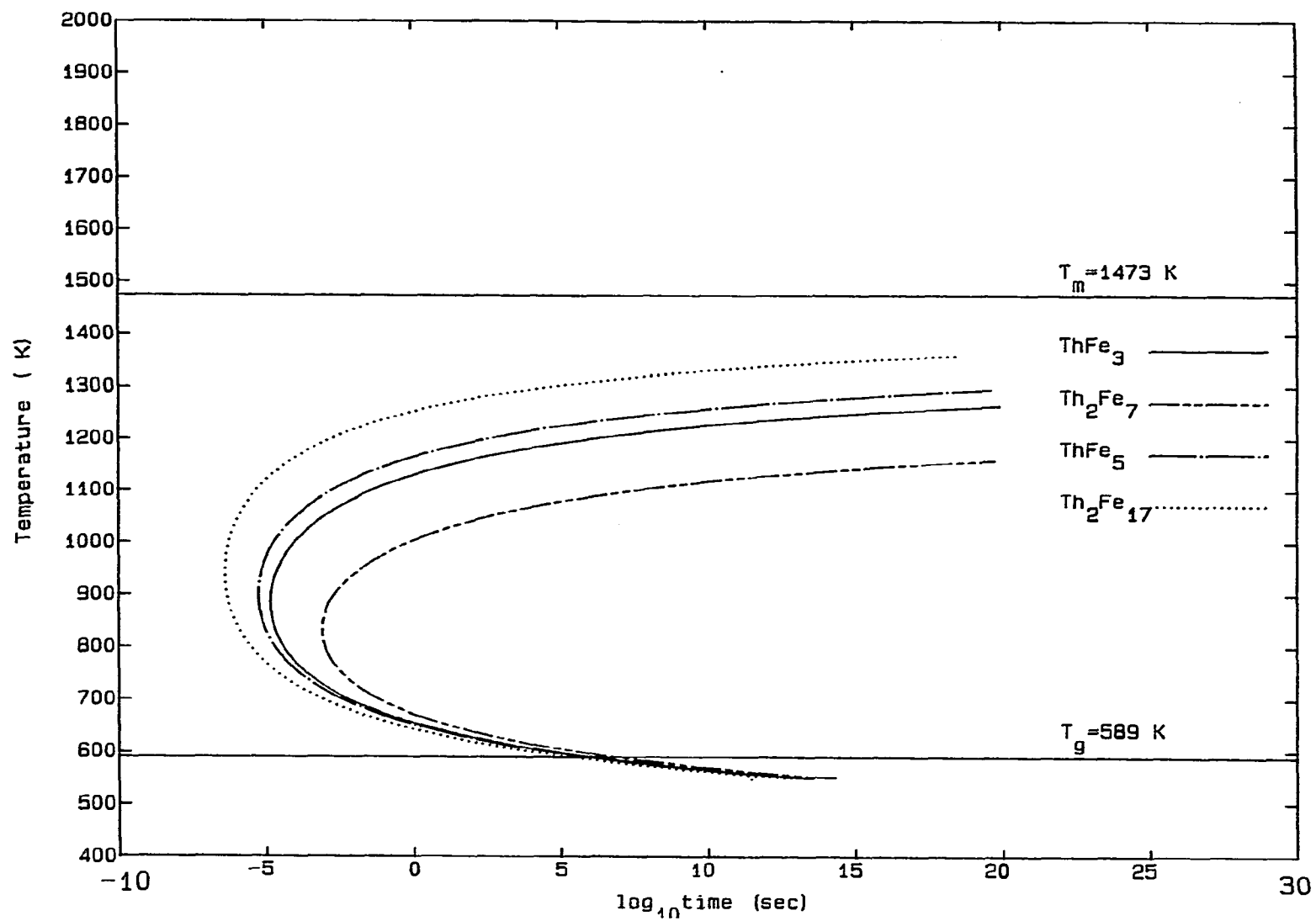


Fig. 12 TTT curves (30 at.% Th)

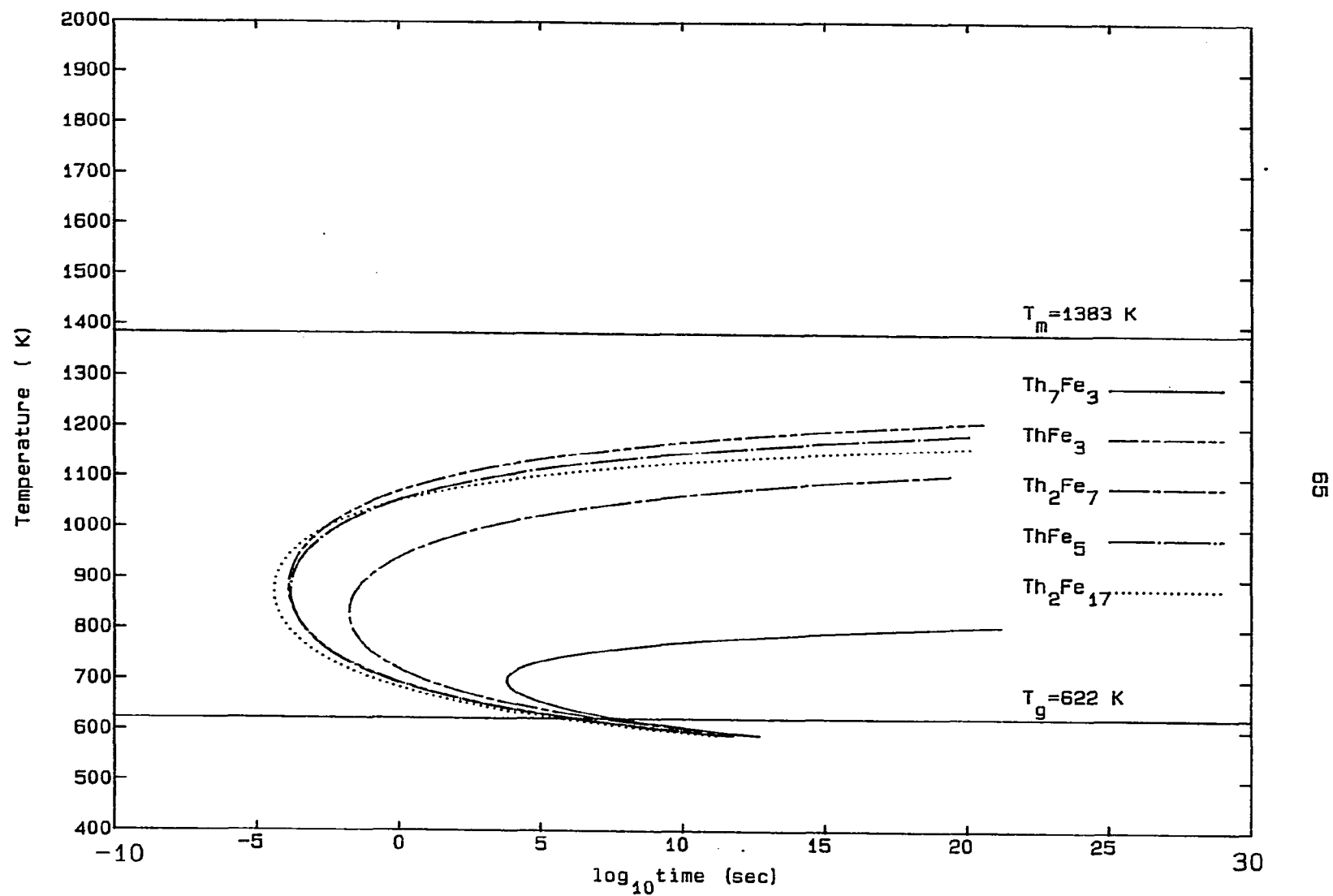


Fig. 13 TTT curves (40 at.% Th)

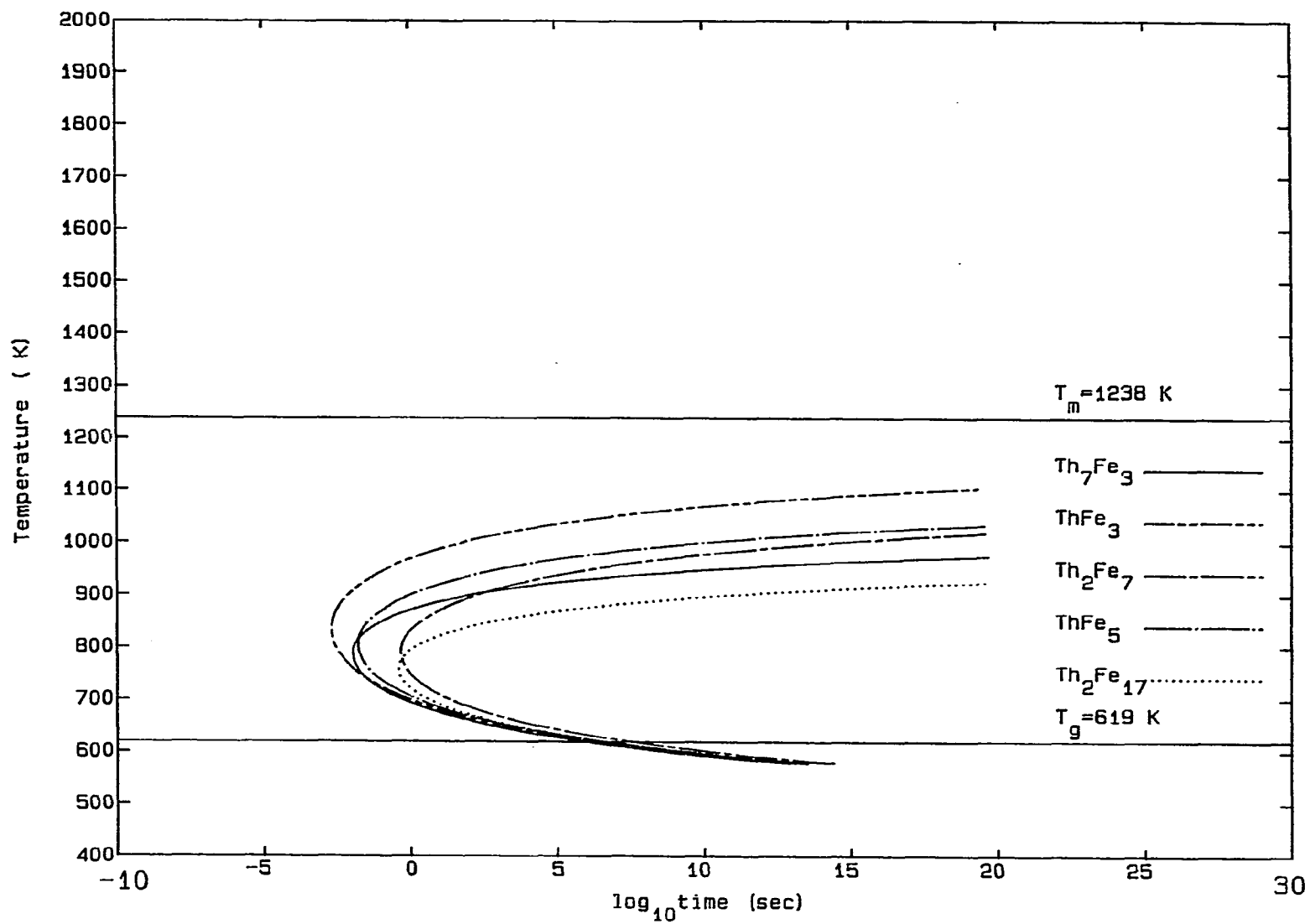


Fig. 14 TTT curves (50 at.% Th)

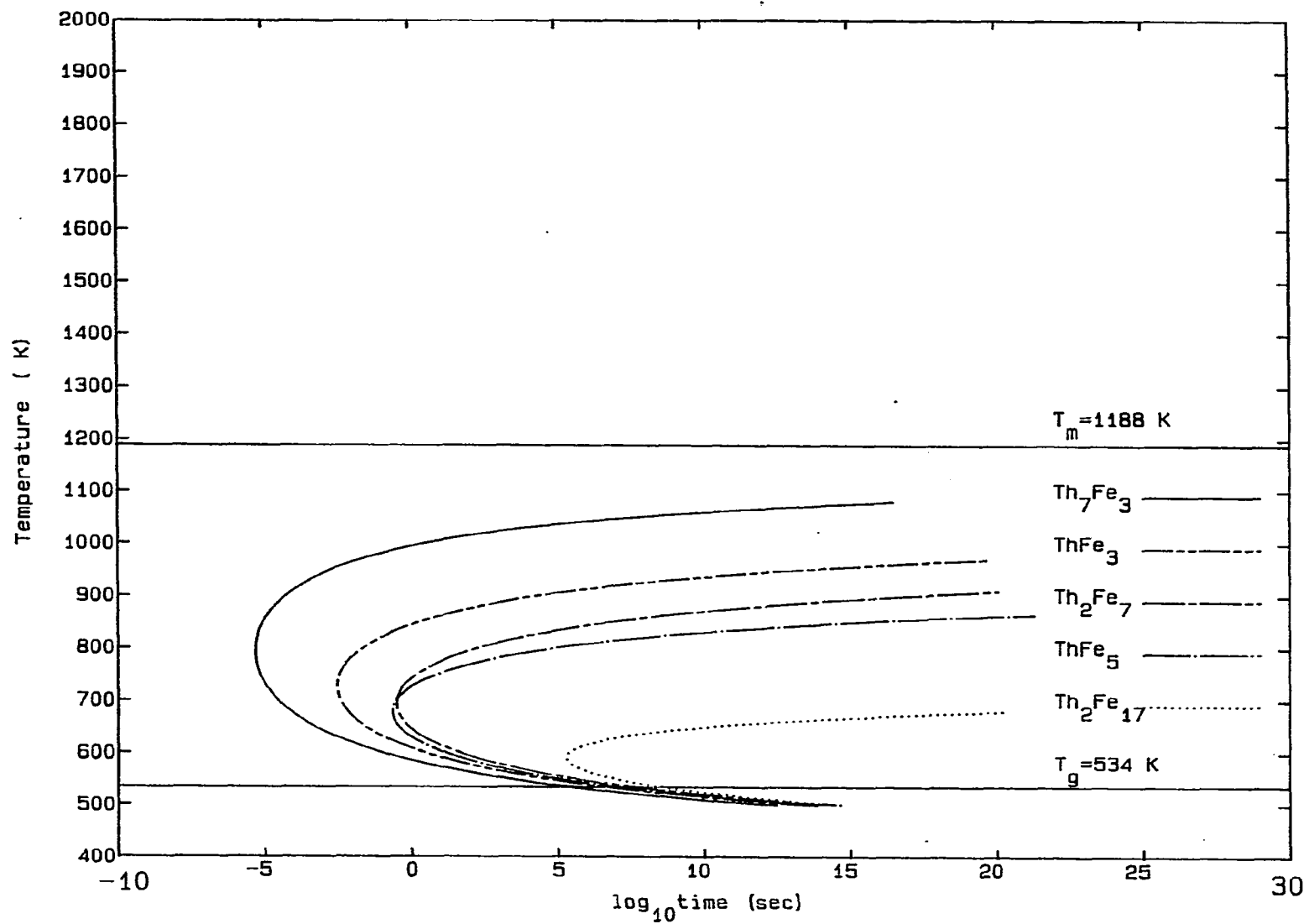


Fig. 15 TTT curves (60 at.% Th)

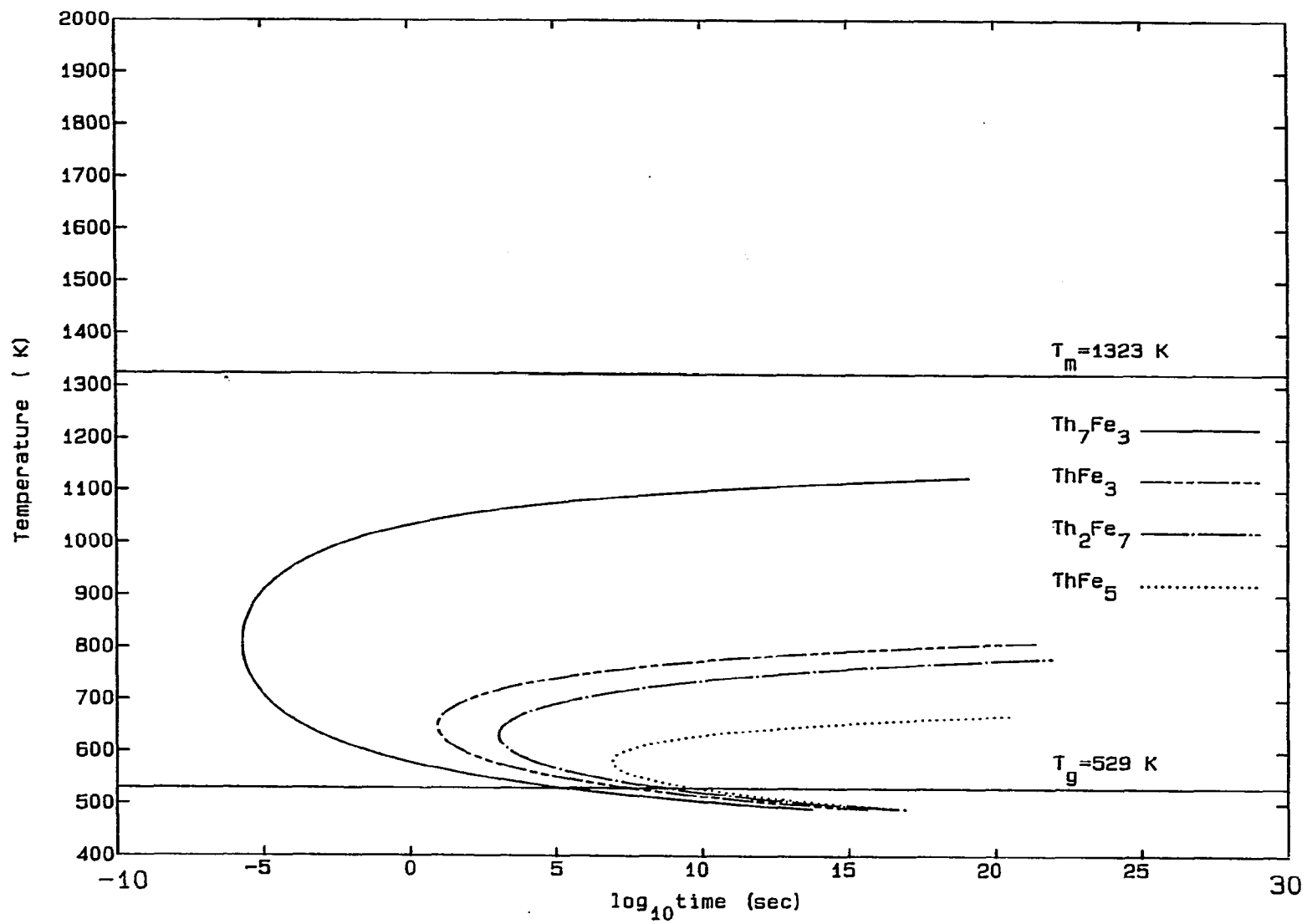


Fig. 16 TTT curves (70 at.% Th)

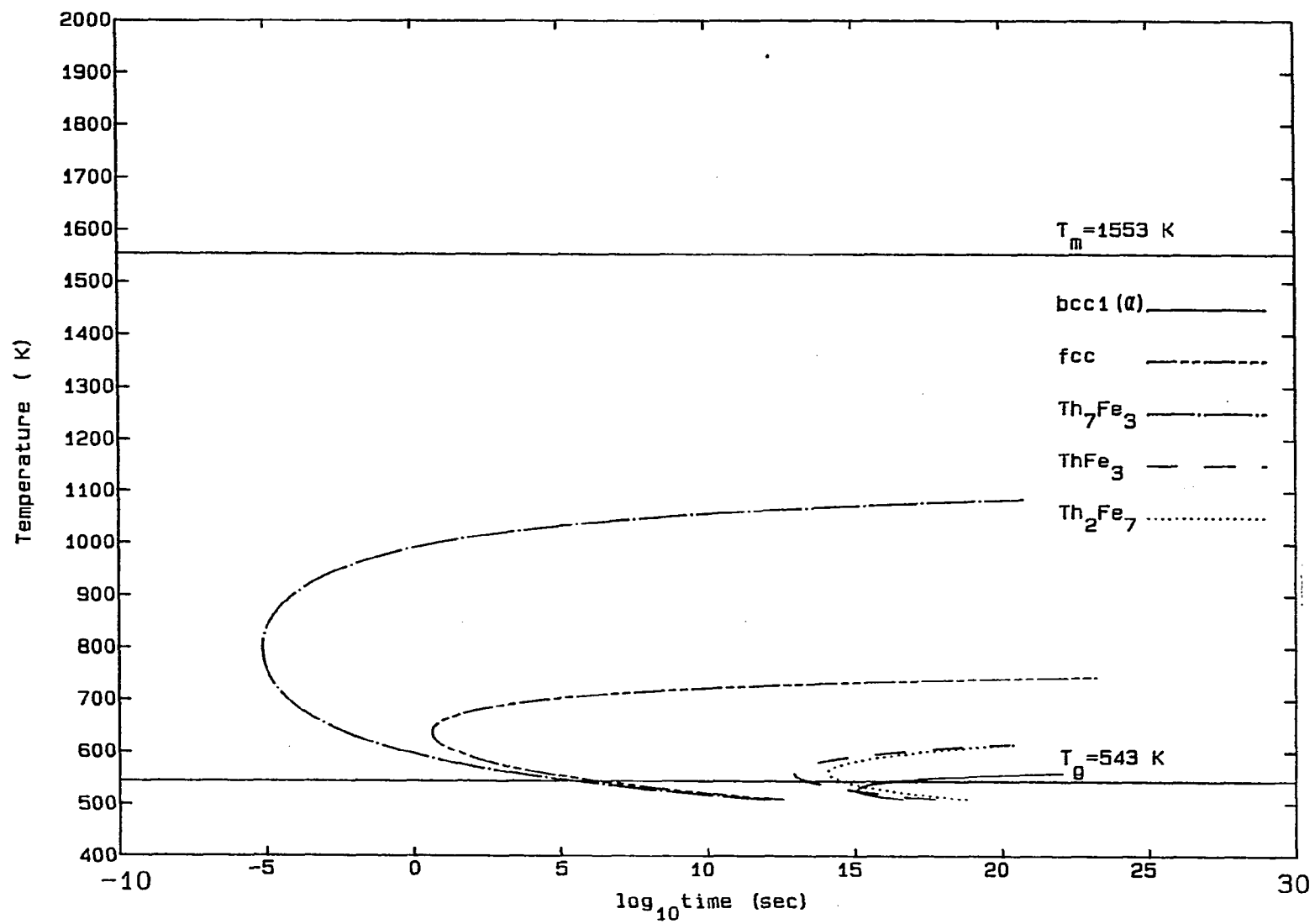


Fig. 17 TTT curves (80 at.% Th)

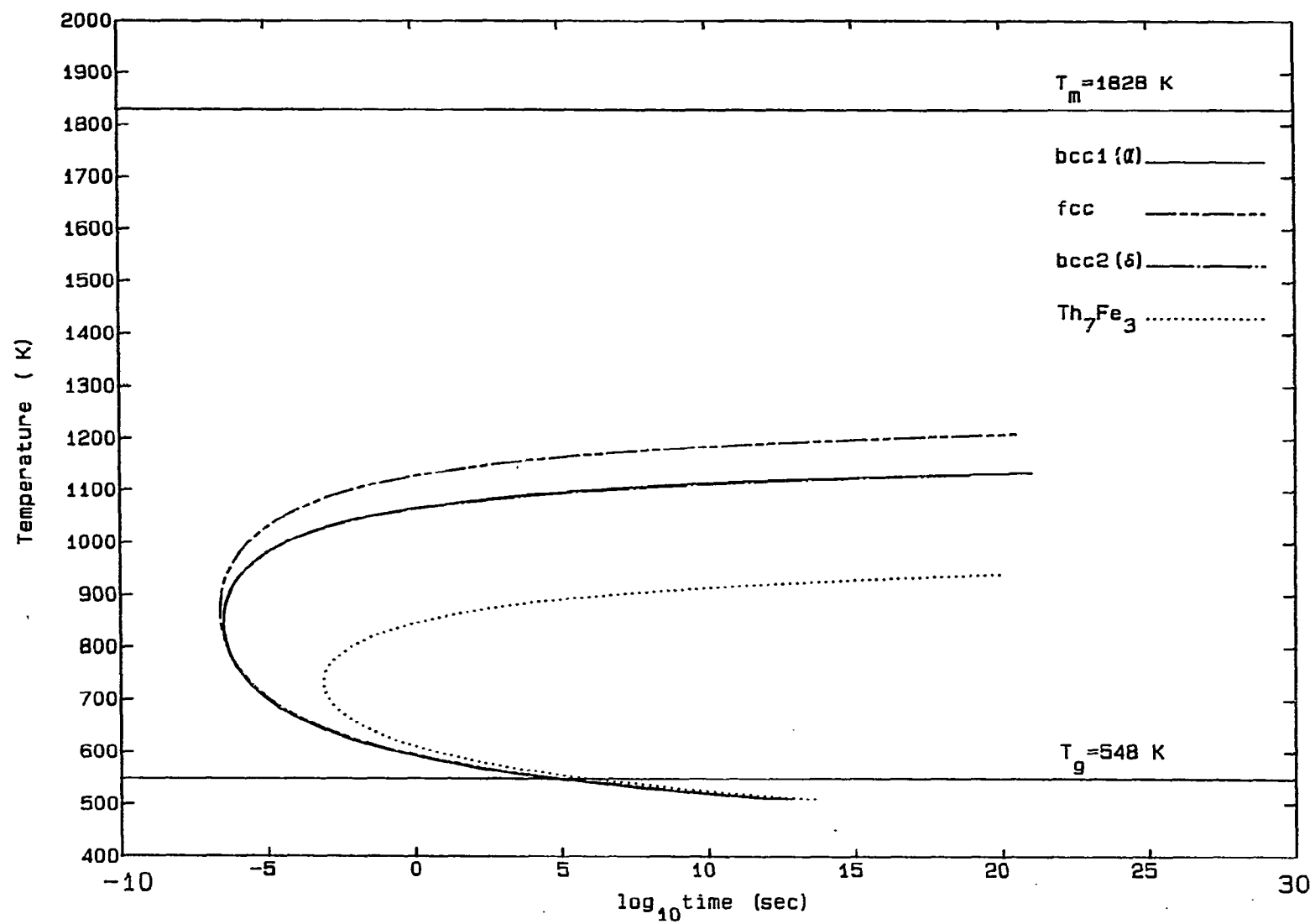


Fig. 18 TTT curves (90 at.% Th)

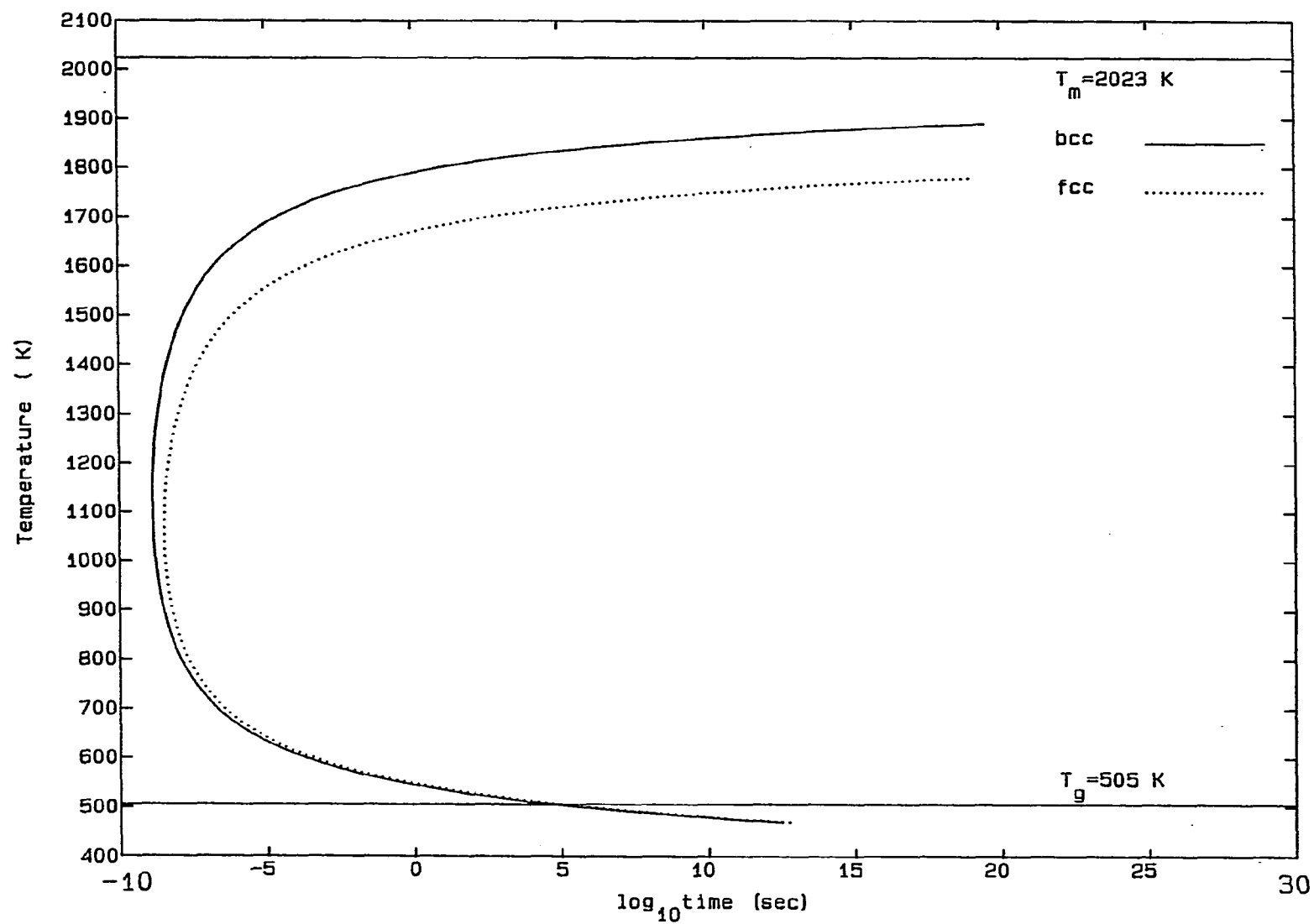


Fig. 19 TTT curves (100 at.% Th)



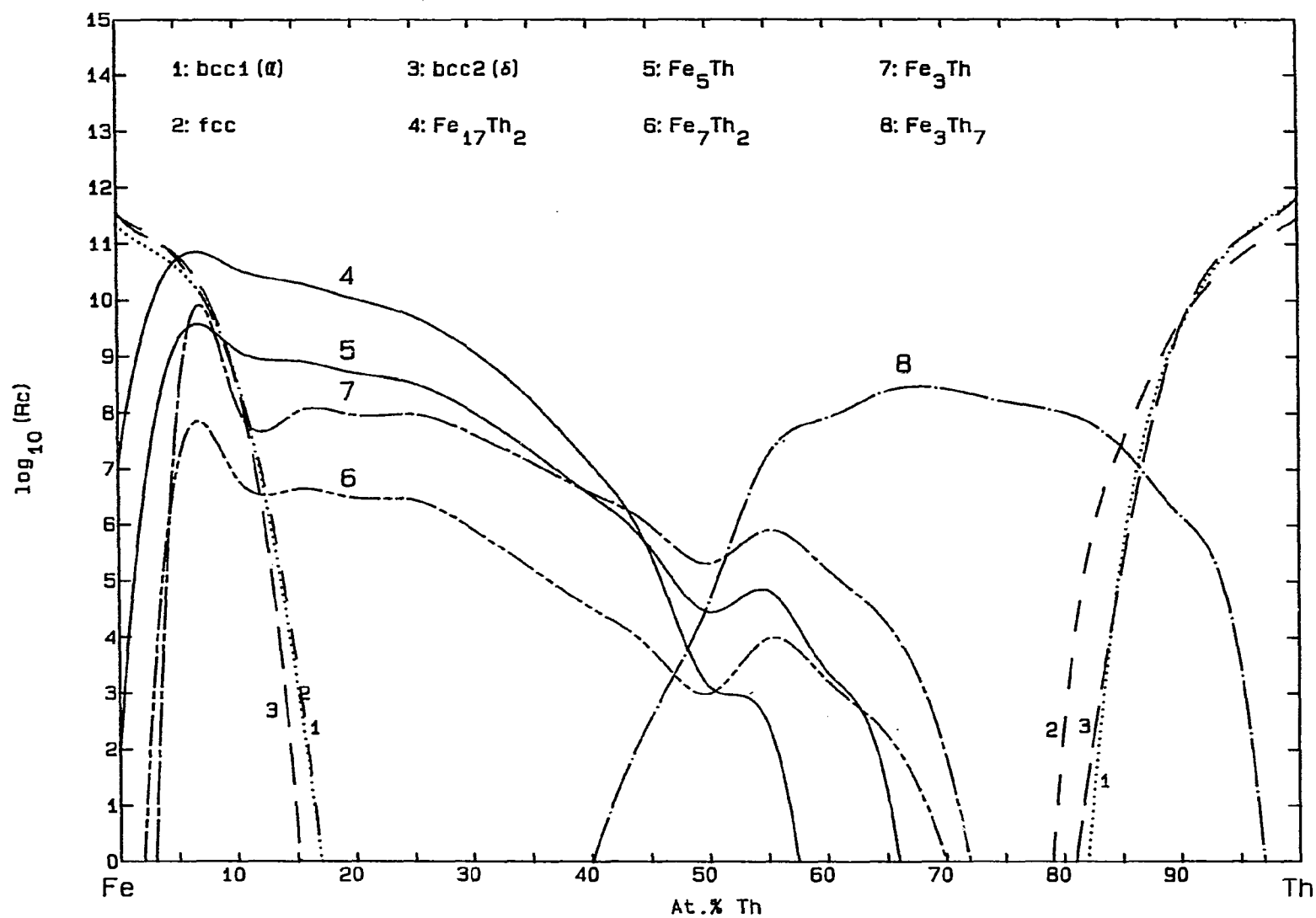


Fig. 20 Critical cooling rate ( $R_c$ ) for solid phases

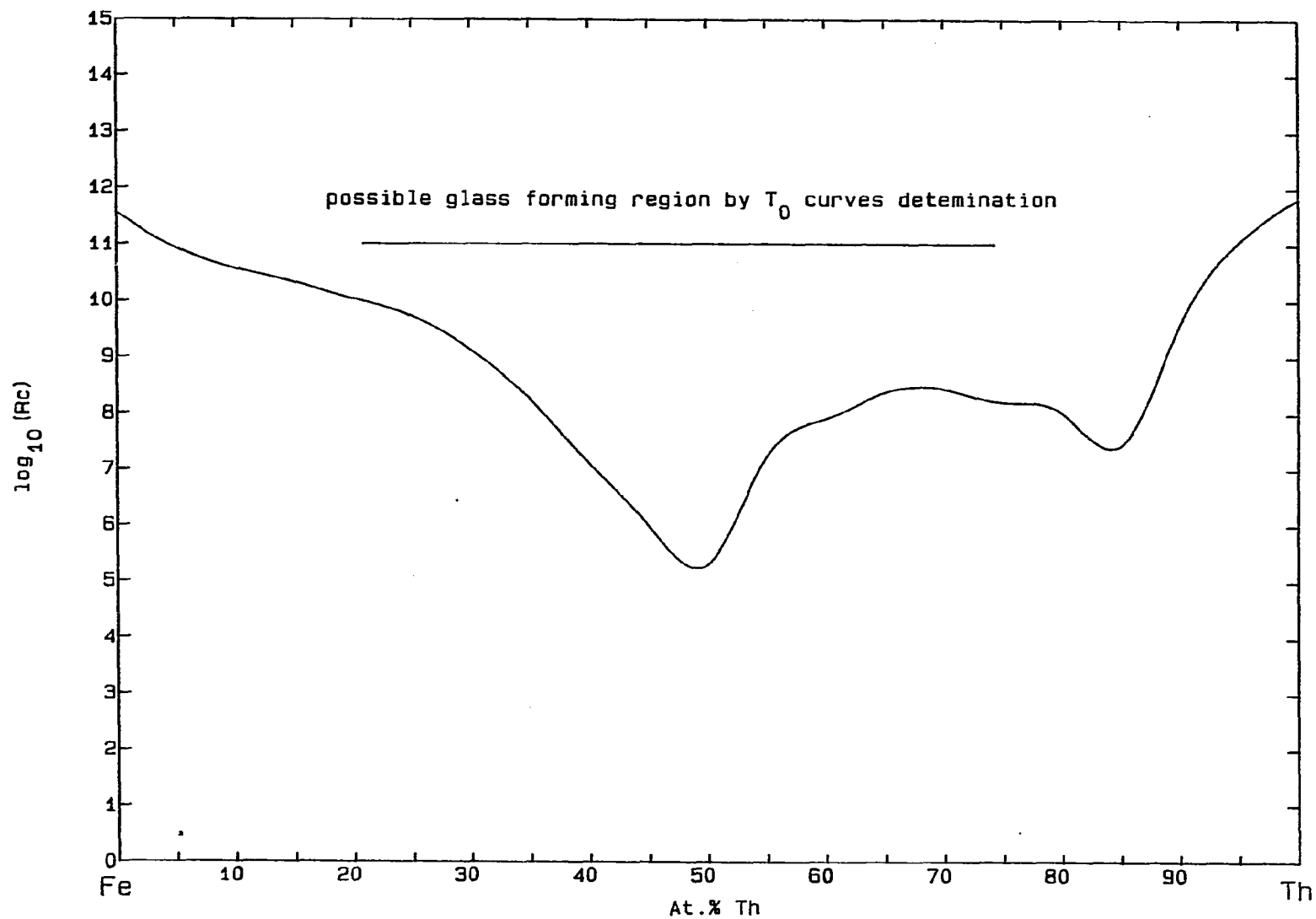


Fig. 21 Critical cooling rate ( $R_c$ ) in the Fe-Th system

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## OVERALL SUMMARY

It was not the purpose of this study to reevaluate the phase diagram, but rather to generate internally consistent thermodynamics functions for each phase. From the phase diagram information, mostly invariant reactions, experimental data and a theoretical model, such thermodynamic descriptions for the phases were obtained with input information. First, all available thermodynamic data were assembled and a model which satisfactorily fits the experimental data was developed. Since any regression creates a best fit to all data, parameters were reevaluated such that they predicted certain critical phase diagram areas exactly along with representing the thermodynamic property. The second and third step in this study was set up in a way that each step optimized the parameters such that the difference between evaluated phase diagram and calculated one is approximately zero.

This analysis indicates that thermodynamic data alone can not be relied upon to predict the phase boundary behavior. Both thermodynamic data and phase diagram data should be used.

The glass forming ability of Fe-Th alloys was determined from TTT and  $T_0$  curves with data for the thermodynamic and kinetic parameters being based upon input

from the calculation of the equilibrium diagram. It has thus been demonstrated that the Gibbs energy functions describing the equilibrium phase diagram can be useful in deriving metastable as well as stable phase formation, and a complete assessment of the metastable equilibria should include consideration of equilibrium thermodynamic functions as well as kinetic factors.

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